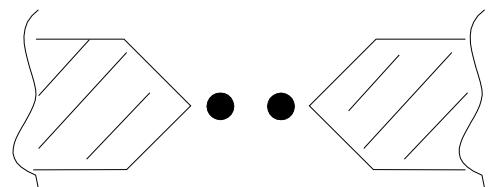
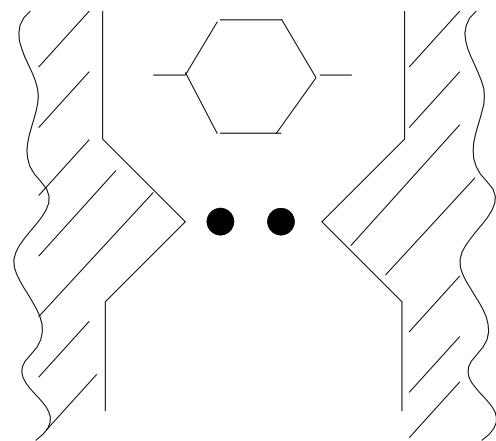


Coherent electron transport, Landauer-Buttiker approach



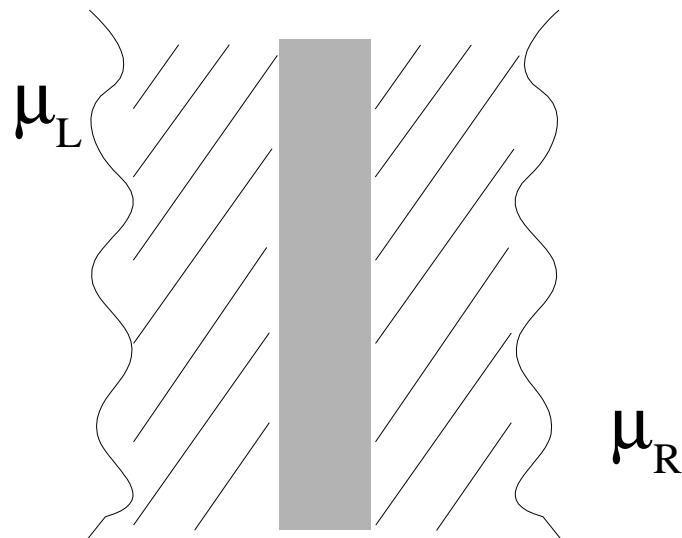
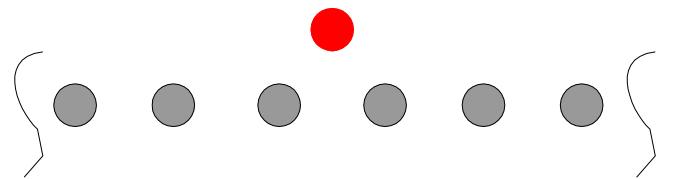
$$\begin{array}{c} \leftarrow \\ L \\ \rightarrow \end{array}$$

$$L < l, L < L_\varphi, L \sim \lambda_F$$

l - mean free path

L_φ - phase decoherence length

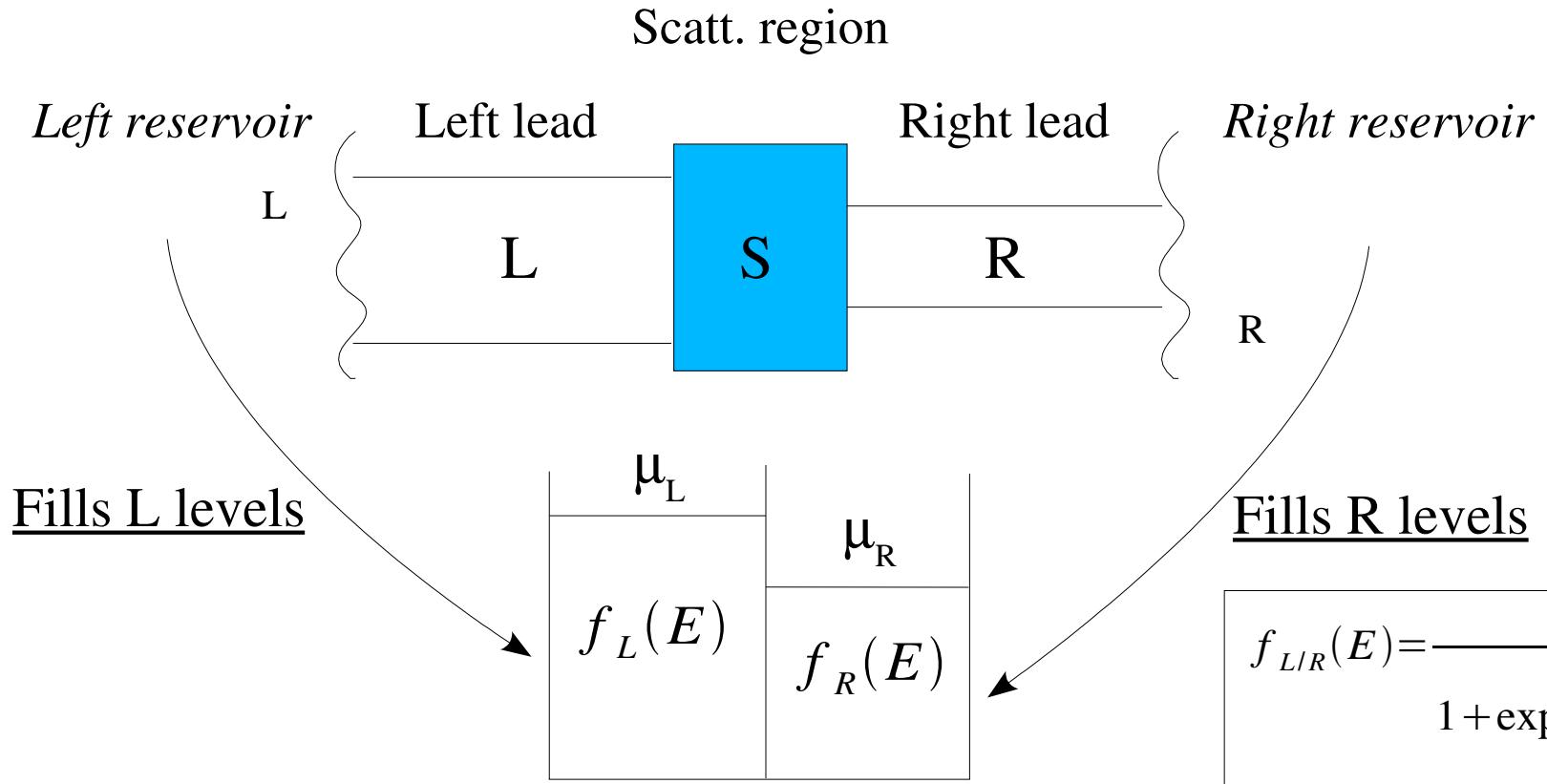
λ_F - wave length



$$\begin{array}{c} I \longrightarrow \\ I = ? \end{array}$$

Coherent ballistic transport, quantum mechanics is applicable, conductance is given by Landauer-Buttiker formula

Landauer-Buttiker formula



$$f_{L/R}(E) = \frac{1}{1 + \exp\left[\frac{E - \mu_{L/R}}{kT}\right]}$$

Landauer-Buttiker formula:

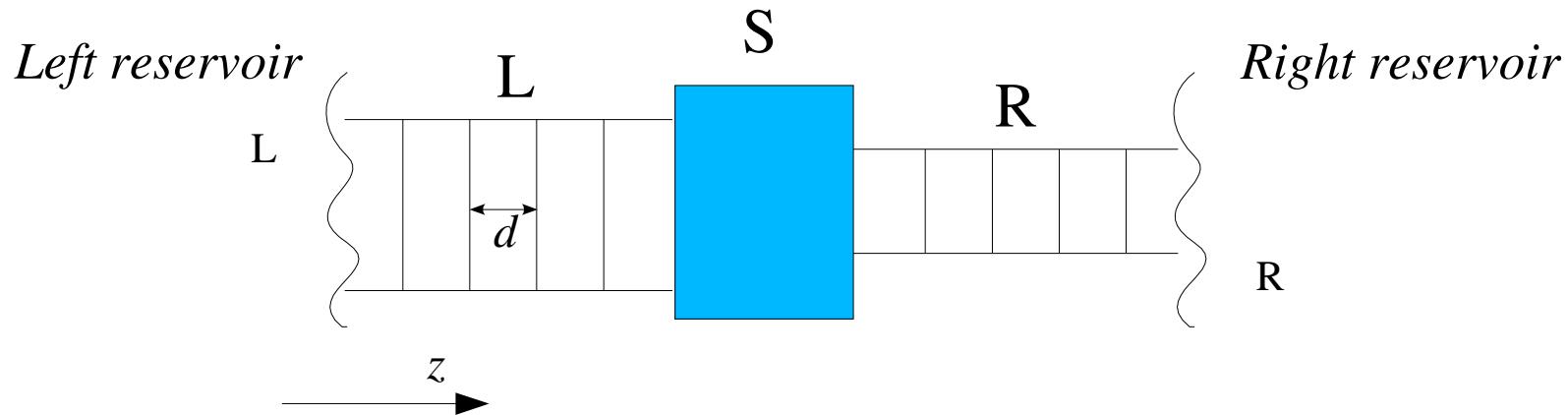
$$I = \frac{e}{h} \int T(E) [f_L(E) - f_R(E)] dE$$

$$\text{Zero temperature, } I = \frac{e}{h} \int_{\mu_R}^{\mu_L} T(E) dE$$

In the linear response regime, at infinitely small voltage, $\mu_L - \mu_R = e \delta V$,

$$G = \frac{I}{\delta V} = \frac{e^2}{h} T(E_F).$$

Asymptotic wave functions in the leads, complex band structure

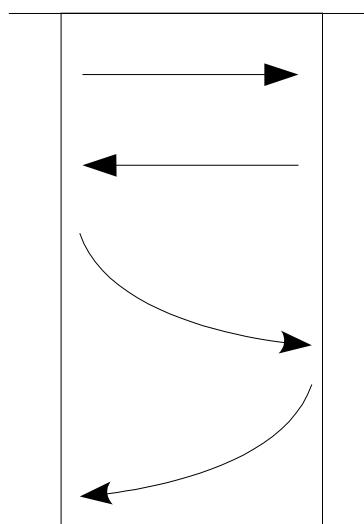


The leads are preperiodic in the direction of transport, the z axis:

$$\varphi(x, y, z) = \exp(ikz)u(x, y, z) \quad - \text{ Bloch form}$$

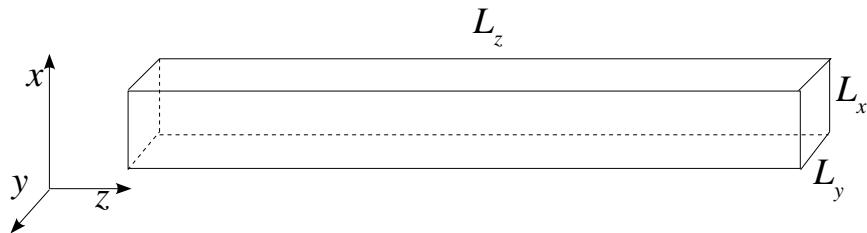
L or R

At fixed energy E :



- $\varphi_{L/R, j}^+$ - propagating to the right, real k , $M_{L/R}(E)$ number of modes
 $\varphi_{L/R, j}^-$ - propagating to the left, real k , $M_{L/R}(E)$ or channels
 $\tilde{\varphi}_{L/R, j}^+$ - decaying to the left, complex k , $\text{Im}k > 0$, Infinite number
 $\tilde{\varphi}_{L/R, j}^-$ - decaying to the right, complex k , $\text{Im}k < 0$, Infinite number

Complex band structure, example: hard wall potential



$$V(x, y, z) = \begin{cases} 0, & 0 < x < L_x, \quad 0 < y < L_y \\ \infty, & \text{otherwise} \end{cases}$$

$$\varphi_{mnk} = \frac{1}{\sqrt{L_z}} \chi_{mn}(x, y) \exp(ikz)$$

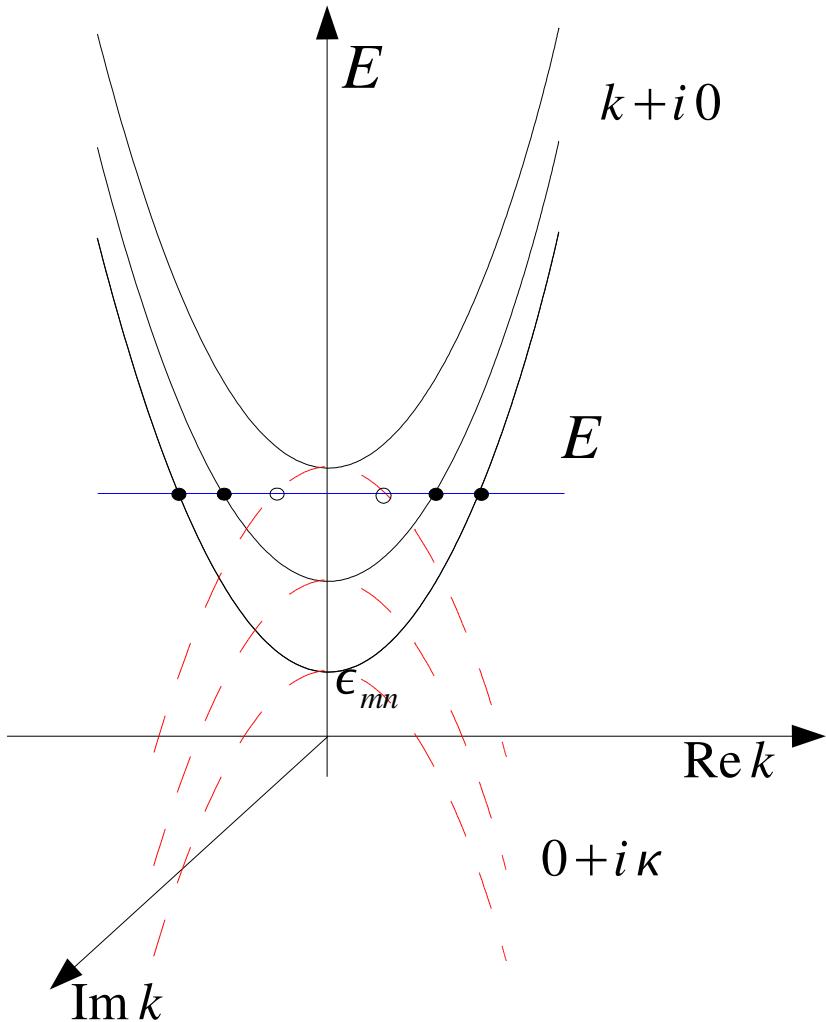
$$\chi_{mn}(x, y) = \frac{1}{\sqrt{L_x L_y}} \sin\left(m \frac{\pi}{L_x} x\right) \sin\left(n \frac{\pi}{L_y} y\right)$$

$$E = \epsilon_{mnk} = \epsilon_{mn} + \frac{\hbar^2 k^2}{2\mu}, \quad \epsilon_{mn} = \frac{\hbar^2 \pi^2}{2\mu L_x^2} m^2 + \frac{\hbar^2 \pi^2}{2\mu L_y^2} n^2$$

\$m, n\$ — integers, numbering transverse subbands

$$k = \sqrt{\frac{2\mu}{\hbar^2} (E - \epsilon_{mn})} \quad \text{and is real or } \textit{imaginary}, \text{ depending on } E,$$

Complex band structure, example: hard wall potential



if $E > \epsilon_{mn}$

$$\varphi_j^{\pm} = \frac{1}{\sqrt{L_z}} \chi_{mn}(x, y) \exp(\pm ikz)$$

$$k = \sqrt{\frac{2\mu}{\hbar^2}(E - \epsilon_{mn})}$$

propagating to the right and to the left modes

Number of modes or channels, $M(E) =$ number of real bands crossing $E =$ number of $\epsilon_{mn} < E$, 2 in our case

if $E < \epsilon_{mn}$

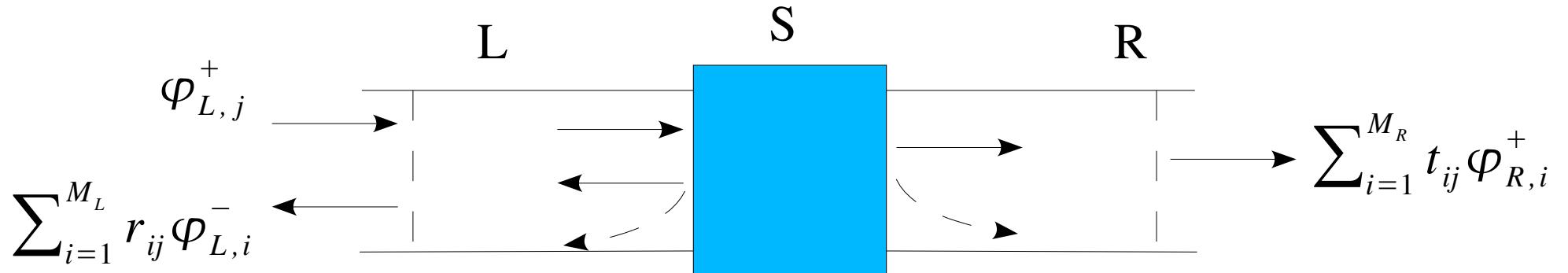
$$\tilde{\varphi}_j^{\pm} = \frac{1}{\sqrt{L_z}} \chi_{mn}(x, y) \exp(\mp \kappa z)$$

$$k = \pm i\kappa$$

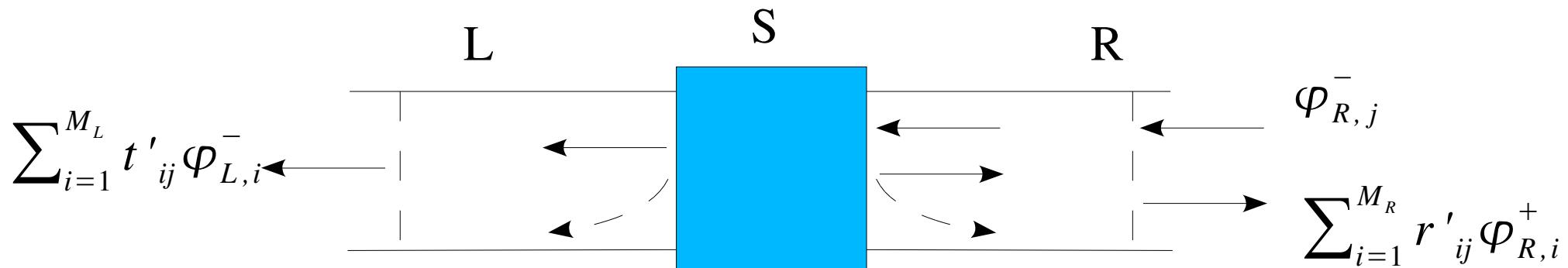
$$\kappa = \sqrt{\frac{2\mu}{\hbar^2}(\epsilon_{mn} - E)}$$

decaying to the right and to the left modes

S-matrix and transmission coefficient



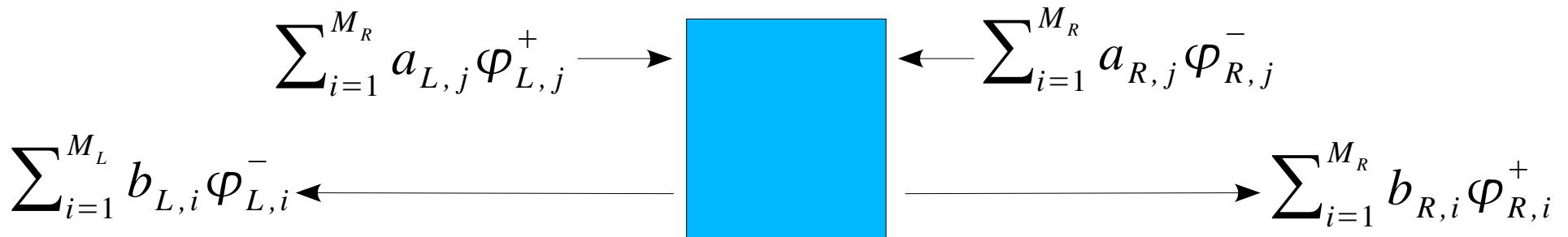
The scattering state $\Phi_{L,j}^+$ originated from $\varphi_{L,j}^+, M_L(E)$



The scattering state $\Phi_{R,j}^-$ originated from $\varphi_{R,j}^-, M_R(E)$

General solution of the SE at the energy E is a linear combination $\left\{ \Phi_{L,j}^+, \Phi_{R,j}^- \right\}$

S-matrix and transmission coefficient (continuation)



$$[a] = \begin{bmatrix} a_L \\ a_R \end{bmatrix} \quad - \text{ given}$$

Related by means of S-matrix:

$$[b] = [S][a]$$

$$[b] = \begin{bmatrix} b_L \\ b_R \end{bmatrix} \quad - \text{ to find}$$

$$[S] = \begin{bmatrix} [r] & [t'] \\ [t] & [r'] \end{bmatrix} \underbrace{M_L + M_R}_{M_L + M_R}$$

Ballistic transport with PWCOND code

Running sequence:

1. Supercell pw.x calculations



- Self-consistent potential for left/right leads and the scattering region

2. pwcond.x calculation



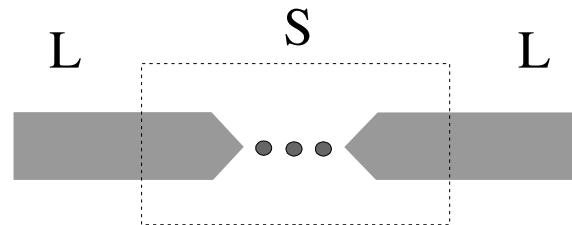
- Complex band structure (CBS) of the leads
- Transmission coefficients for each propagating channel of the left lead

Three kinds of calculations controlled by IKIND flag

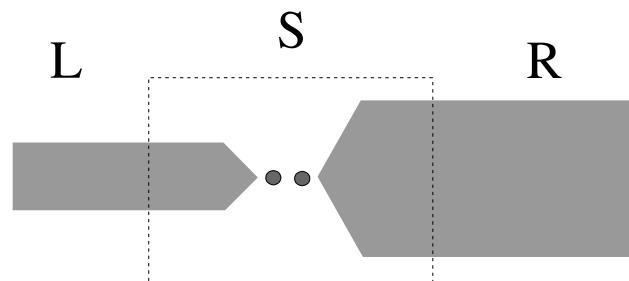
1. IKIND = 0 :

CBS of the lead

2. IKIND = 1 : The scattering problem with *identical* leads



3. IKIND = 2 : The scattering problem with *different* leads

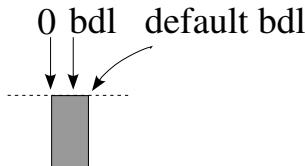


Separate files for LEADS and SCATTERING region

- IKIND = 0

lead

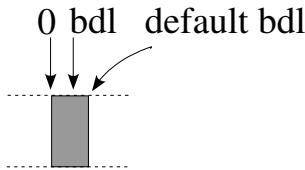
plefixl = '...'



- IKIND = 1

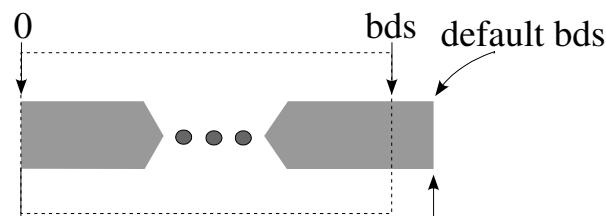
left (right) lead

plefixl = '...'



scattering
region

plefixs = '...'

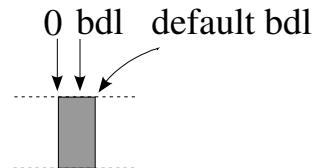


unit cell for pw.x

- IKIND = 2

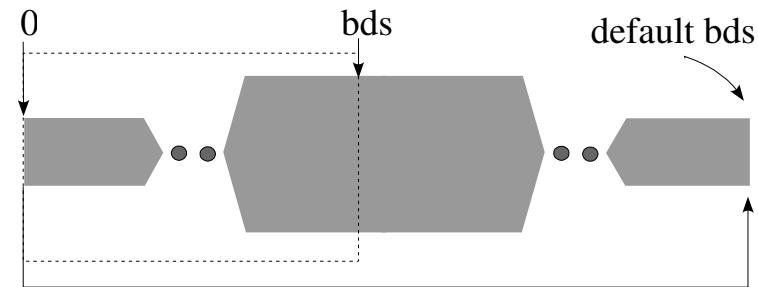
left lead

plefixl = '...'



scattering
region

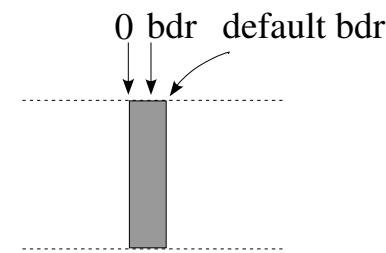
plefixs = '...'



unit cell for pw.x

right lead

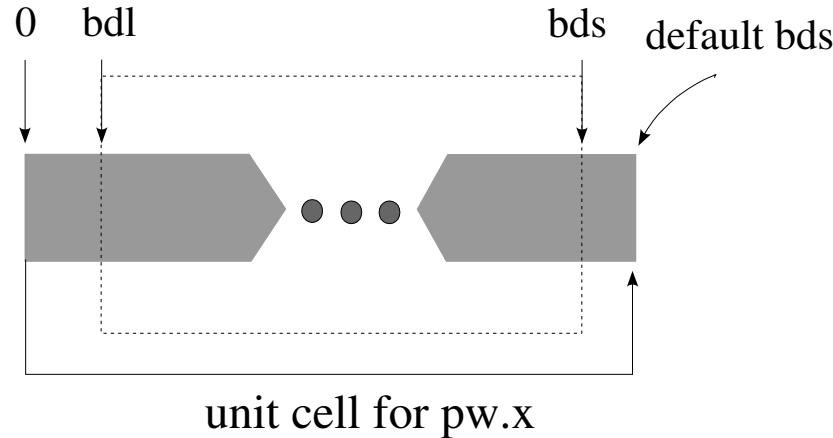
plefixr = '...'



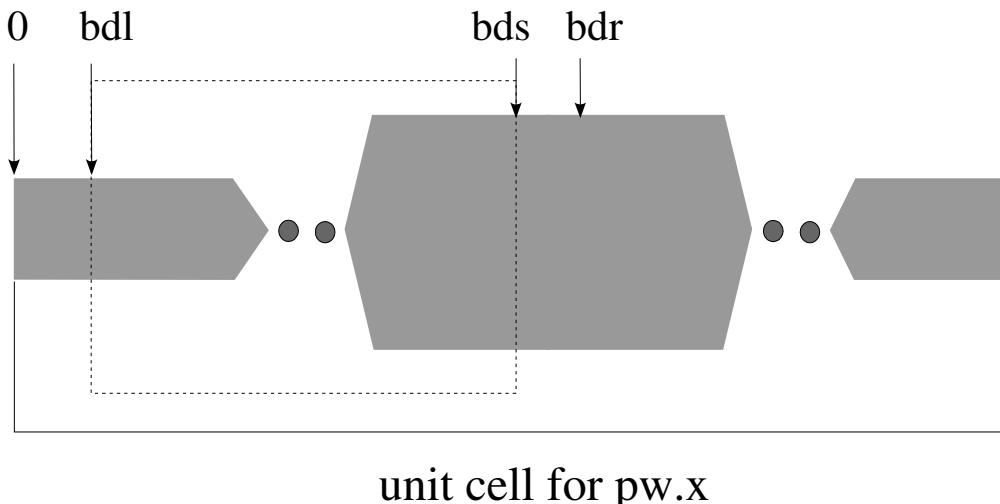
A single file containing all the regions

plefixt = '...'

- IKIND = 1



-
- IKIND = 2



Input variables of PWCOND (see Doc/INPUT_PWCOND.txt)

&inputcond

Already discussed

ikind,
prefixl, prefixs, prefixr, prefixt,
bdl, bds, bdr

Controlling the accuracy:

ewind, Larger ewind and nz1, and smaller epsproj
epsproj, ↓
nz1 higher accuracy

Output of CBS and transmission

band_file,
tran_file

Mesh in energy

energy0 - initial energy (in eV from Ef)
denergy - step in energy (in eV)

/

| | | |
|----------------------|--------------------|--------------------------|
| <code>nkpts</code> | - | number of (kx,ky) points |
| <code>kx(1)</code> | <code>ky(1)</code> | <code>weight(1)</code> |
| <code>kx(2)</code> | <code>ky(2)</code> | <code>weight(2)</code> |
| ... | | |
| <code>nenergy</code> | - | number of energies |

If `denergy = 0` the list of energies is provided after `nenergy`

References

- *Books on electron transport:*
S. Datta “Electronic transport in mesoscopic systems”;
D.K. Ferry and S.M. Goodnick “Transport in nanostructures”

- *Details of the method implemented in PWCOND:*
H.J. Choi and J. Ihm, Phys. Rev. B 59, 2267 (1999);
A. Smogunov, A. Dal Corso, and E. Tosatti, Phys. Rev. B 70, 045417 (2004)