





Imaging quantum turbulence in ³He-B: Do spectral properties of Andreev reflection reveal properties of turbulence?

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To validate the experimental visualization of turbulence in ³He-B by showing the relation between the vortex-line density and the Andreev reflectance of the vortex tangle in the first simulations of the Andreev reflectance by a realistic 3D vortex tangle, and comparing the results with the first experimental measurements in ³He-B

able to probe quantum turbulence on length scales smaller than the inter-vortex separation¹.

¹In Prague ⁴He experiments at temperaturs $1 \text{ K} < T < T_{\lambda}$, quantum turbulence on the length scales smaller than the intervortex separation was visualized by the Lagrangian dynamics of micrometer-sized solid particles, see e.g. La Mantia and Skrbek, PRB **90**, 014519 (2014).

Superfluidity and quantized vortices in ⁴He and ³He-B



Transition temperature

 $T_{\lambda} \approx 2.17 \,\mathrm{K}$

 $T_c\approx 1\,\mathrm{mK}$

Quantum of circulation

$$\kappa = rac{2\pi\hbar}{m_4} pprox 0.997 imes 10^{-7} \, rac{\mathrm{m}^2}{\mathrm{s}} \hspace{1.5cm} \kappa = rac{2\pi\hbar}{2m_3} pprox 0.662 imes 10^{-7} \, rac{\mathrm{m}^2}{\mathrm{s}}$$

We will be interested in dynamics of quantized vortex lines and associated flow field at temperatures $\ll T_c$ so that the normal fluid can be neglected.

Quantum turbulence in the zero temperature limit

Each vortex moves in a $\underline{\text{collective flow field}}$ of all other vortices

 \implies complex dynamics, accompanied by vortex reconnections

⇒ vortex tangle (quantum turbulence)

Regimes and spectra of quantum turbulence

Quasiclassical

Ultraquantum



Intencity of quantum turbulence is characterized by the vortex line density, $L \ [m^{-2}]$

- total length of vortex lines

per unit volume



Large-scale flow



No large-scale flow. The only length scale is the intervortex distance.

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Detection and visualization of quantum turbulence

| $\frac{^{4}\text{He:}}{1 \text{ K}} 1 \text{ K} \lesssim \mathcal{T} \leq \mathcal{T}_{\lambda} \approx 2.17 \text{ K}$ | <u>³He-B:</u> N/A for $0.25 T_c < T < T_c$ |
|--------------------------------------------------------------------------------------------------------------------------|-------------------------------------------------------------------------------------------------|
| Second sound | $\frac{^{3}\text{He-B:}}{T}$ $T < 0.25 T_{c}$ |
| PIV and PTV (Particle Tracking) | Andreev reflection of quasiparticle thermal excitations |
| He_2^* excimer molecules | (not yet a truly visualization technique) |
| <u>⁴He:</u> <i>T</i> < 1 K | Advantages: 1) Non-invasive |
| He^*_2 excimer molecules | 2) in combination with numerical simulations, can |
| Electron bubbles | (in particular, allows to measure and interpret <u>fluctuations</u> of the vortex line density) |
| | Fisher et al. , PRL 63 (1983); Fisher, Jackson, Sergeev, Tsepelin, PNAS 111 (2014) |

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Governing equations

$$\mathbf{v}(\mathbf{r}, t) = -rac{\kappa}{4\pi} \oint_{\mathcal{L}} rac{\mathbf{r} - \mathbf{s}}{|\mathbf{r} - \mathbf{s}|^3} imes d\mathbf{s}, \quad rac{d\mathbf{s}}{dt} = \mathbf{v}(\mathbf{s}, t)$$

 solved in the cubic box of size 1 mm using the vortex-filament method with periodic boundary conditions. Biot-Savart integrals are evaluated via a tree method.

[Baggaley & Barenghi, PRB 84, 020504 (2011)]



Tangle is generated by vortex loops (radius 200 μ m) injected at a frequency $f_{\rm f} = 10$ Hz. To insure isotropy, the loop injection plane is switched at both corners at frequency $f_{\rm f} = 3.3$ Hz.

Evolution of the vortex line length Saturated, time-averaged line density in the statistically steady state: $\langle L \rangle = 9.7 \times 10^7 \,\mathrm{m^{-2}}$.



[Baggaley, Tsepelin, Barenghi, Fisher, Pickett, Sergeev, and Suramlishvili, PRL 115, 015302 (2015)

Note: the simulated tangle is quasiclassical



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Thermal quasiparticle excitations. And reev scattering

At $T \ll T_c$ excitation propagate ballistically (mean free path \gg size of experimental cell)

Quiescent fluid

$$E = \sqrt{\epsilon_{
ho}^2 + \Delta^2}, \quad {
m where} \quad \epsilon_{
ho} = rac{p^2}{2m^*} - \epsilon_F$$

 Δ – superfluid energy gap; ϵ_F – Fermi energy; p_F – Fermi momentum



- Quasiparticles: $\epsilon_p > 0$
- Quasiholes: $\epsilon_p < 0$

Flowing ³He-B (e.g. quantized vortex)

 $E(\mathbf{p}) \rightarrow E(\mathbf{p}) + \mathbf{p} \cdot \mathbf{v}$ (Galilean transformation)



[effective potential barrier: $U_{\rm eff} = \Delta + \mathbf{p} \cdot \mathbf{v}(\mathbf{r})$] Fisher et al., PRL 63 (1983); Fisher, Jackson, Sergeev, Tsepelin, PNAS 111 (2014)

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Any momentum transfer between excitations and superfluid is minimal

⇒ the reflected excitations almost exactly retrace the path of incoming quasiparticles

Experimental arrangements (Lancaster)

Sketch of experiment ...



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Experimental arrangements (Lancaster)

Sketch of experiment ...



... and its realization



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Andreev reflection of quasiparticle excitations by the vortex tangle

Quasiparticle flux incident in the x-direction on the (y, z)-side of the computational cell:

$$\langle nv_g \rangle_{(y,z)}^i = \int_{\Delta}^{\infty} g(E) f(E) v_g(E) dE$$

[g(E) - density of states, f(E) - Fermi distribution function]

Quasiparticle flux transmitted through the tangle:

$$\langle nv_g \rangle_{(y,z)}^t = g(p_F) \int_{\Delta + p_F V_y^{max}}^{\infty} \exp(-E/k_B T) dE$$

 $[g(p_F) - \text{density of momentum states at the Fermi energy}]$

 V_x^{\max} – the highest superfluid velocity encountered along the quasiparticle's trajectory



The fraction of quasiparticles Andreev-reflected by a tangle along x-direction at position (y, z):

$$f_{y,z} = 1 - \exp\left(-rac{p_F V_x^{\max}}{k_B T}
ight)$$

A 2D map of Andreev reflection by the tangle Baggaley, Tsepelin, Barenghi, Fisher, Pickett, Sergeev, Suramlishvili, PRI 115 015302 (2015)1



(A) Quasiparticle reflection (B) Quasihole reflection

The extended regions of high reflectivity (darker) and low reflectivity (lighter) illustrate the distribution of the large-scale flows

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Total Andreev reflection

Total reflection $f_R = \langle f_{y,z} \rangle$ (averaged over (y, z)-plane)

The results for quasiparticles and quasiholes are combined to yield the reflection for a full thermal beam



Baggaley, Tsepelin, Barenghi, Fisher, Pickett, Sergeev, Suramlishvili, PRL 115, 015302 (2015)

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Spectral properties of fluctuations: Andreev reflection vs quantum turbulence

Statistically steady state

Equilibrium, time-averaged values: $\langle L \rangle = 9.7 \times 10^7 \, \mathrm{m}^{-2}, \qquad \langle f_R \rangle = 0.37$ Fluctuations: $\delta L(t) = L(t) - \langle L \rangle, \ \delta f_R(t) = f_R(t) - \langle f_R \rangle$

Power spectral densities (PSD) of the Andreev reflection for numerical simulations and experimental observations

PSD of the simulated Andreev reflection and the simulated line density



• f^{-3} scaling corresponds to the Andreev reflection on length scales *smaller* than the intervortex distance

• high f – line density spectrum dominated by *unpolarized*, random vortex lines \implies $f^{-5/3}$

• *medium* f – line density spectrum dominated by large scale flows generated by *polarized* vortex lines $\implies f^{-3}$

• cross-over at frequency $f_{\ell} \approx \kappa/(2\pi\ell) \approx 1 \, \mathrm{Hz}$ corresponding to the intervortex distance

• The fluctuations of the vortex line density and of the Andreev reflection are well correlated. However, their spectral densities behave differently: for large scale flows the line density spectrum scales as f^{-3} while that of Andreev reflection as $f^{-5/3}$, for an unpolarized tangle the scalings are reversed.

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Baggaley, Tsepelin, Barenghi, Fisher, Pickett, Sergeev, Suramlishvili, PRL 115, 015302 (2015)

10⁻¹ Frequency. H:

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• The Andreev reflectance of a vortex tangle does indeed reveal the nature of quantum turbulence.

• The $f^{-5/3}$ scaling of the frequency spectrum of the Andreev-retro-reflected signal has been observed earlier in the experiments of the Lancaster ULT group [e.g. Bradley *et al.*, PRL **101**, 065302 (2008)] where it was argued that fluctuations of the reflected signal can be interpreted as fluctuations of the vortex line density.

• The Andreev reflection technique has great potential for elucidating the behavior of pure quantum turbulence.