

Singularities and Irreversibility



Main SPONSOR:

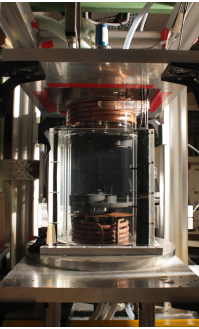
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ANR EXPLOIT: T. Chaabo, A. Cheminet, C. Cuvier, F. Daviaud, P. Debue, B. Dubrulle, D. Geneste-J-M. Foucaut, J-P. Laval, Y. Ostovan, V. Padilla, V. Valori, C. Wiertel

F. Nguyen, H. Faller, D. Kuzzay, D. Faranda, E-W. Saw
C. Nore, J-L. Guermond, L. Cappanera

Symmetries and Navier-Stokes



$$\vec{\nabla} \cdot \vec{u} = 0$$

$$\partial_t \vec{u} + (\vec{u} \cdot \vec{\nabla}) \vec{u} = -\frac{1}{\rho} \vec{\nabla} p + \nu \Delta \vec{u}$$

Broken Symmetry (ii)

Time-reversal

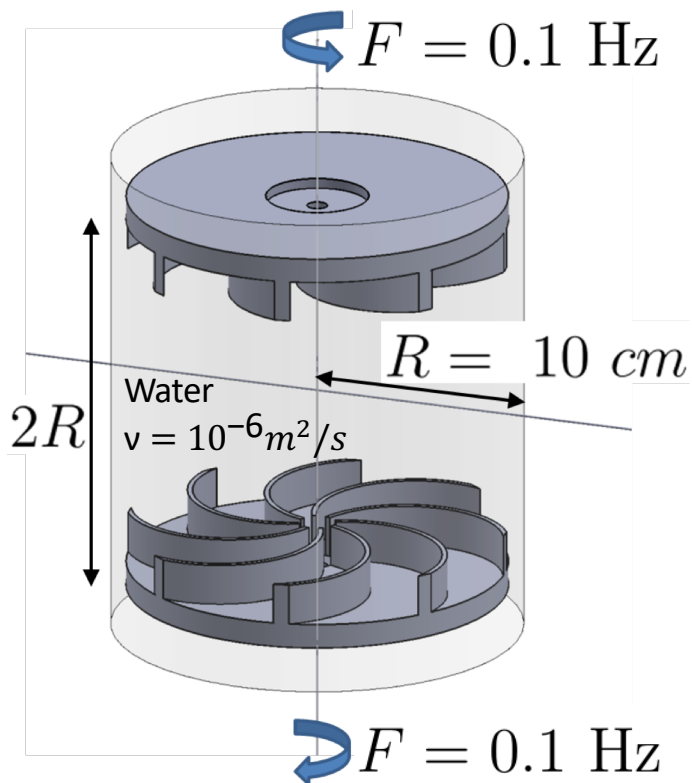
$$t \rightarrow -t$$

Only for $\nu = 0$

$$\mathbf{u} \rightarrow -\mathbf{u}$$

Time-reversal breaking by viscosity

Experimental set-up : the von Kármán flow



- Reynolds number:

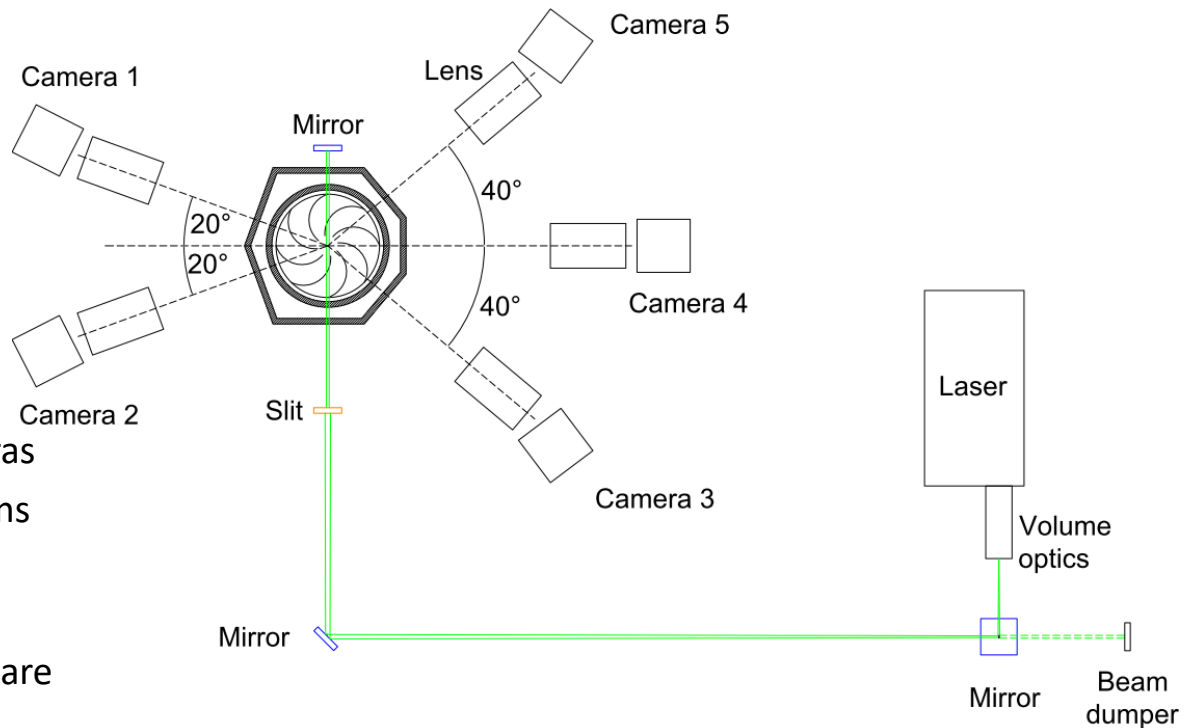
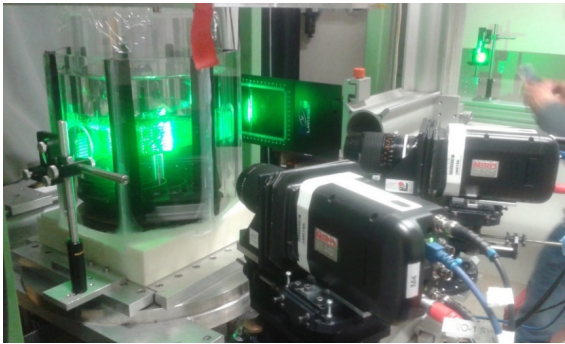
$$Re = \frac{2\pi R^2 F}{\nu} = 6.3 \times 10^3$$

- Average dissipation rate ϵ computed from torque measurements

- Kolmogorov scale :

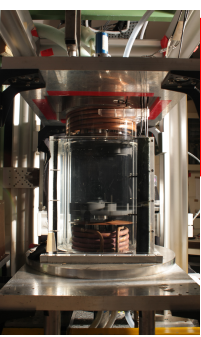
$$\eta = (\nu^3 / \epsilon)^{1/4} = 0.3 \text{ mm}$$

Experimental set-up : velocity measurement

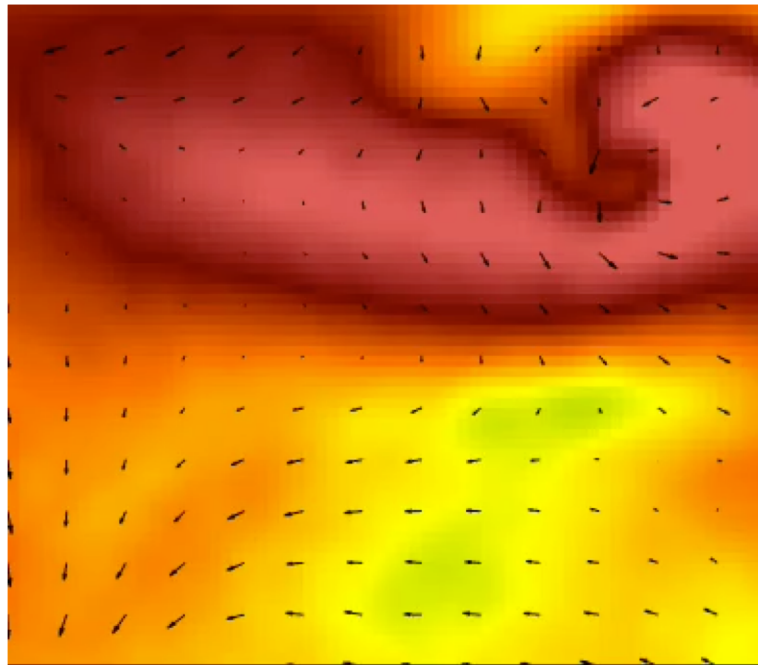


- Tomographic PIV with 5 cameras
- Outer tank to reduce distortions
- Measurement volume :
4cmx4cmx0,5cm
- Processing with Lavisoin software
- 75% overlap
- Vector spacing : $dx=0,35\text{mm}=1,2\eta$
- 30000 frames collected

Velocity measurements in von Karman flow



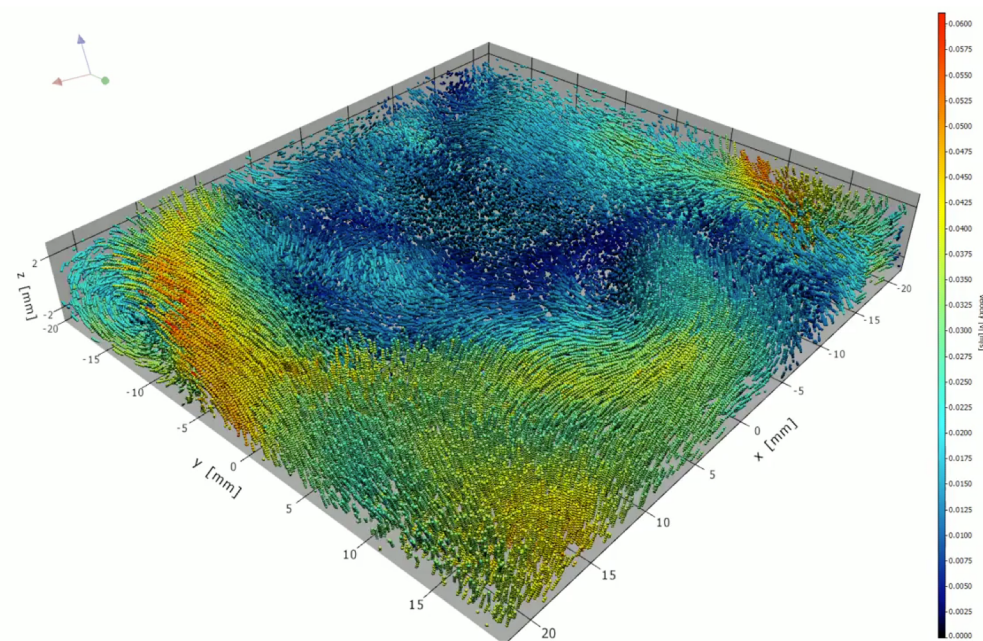
Turbulent flow viewed in PIV



TPIV velocity fields

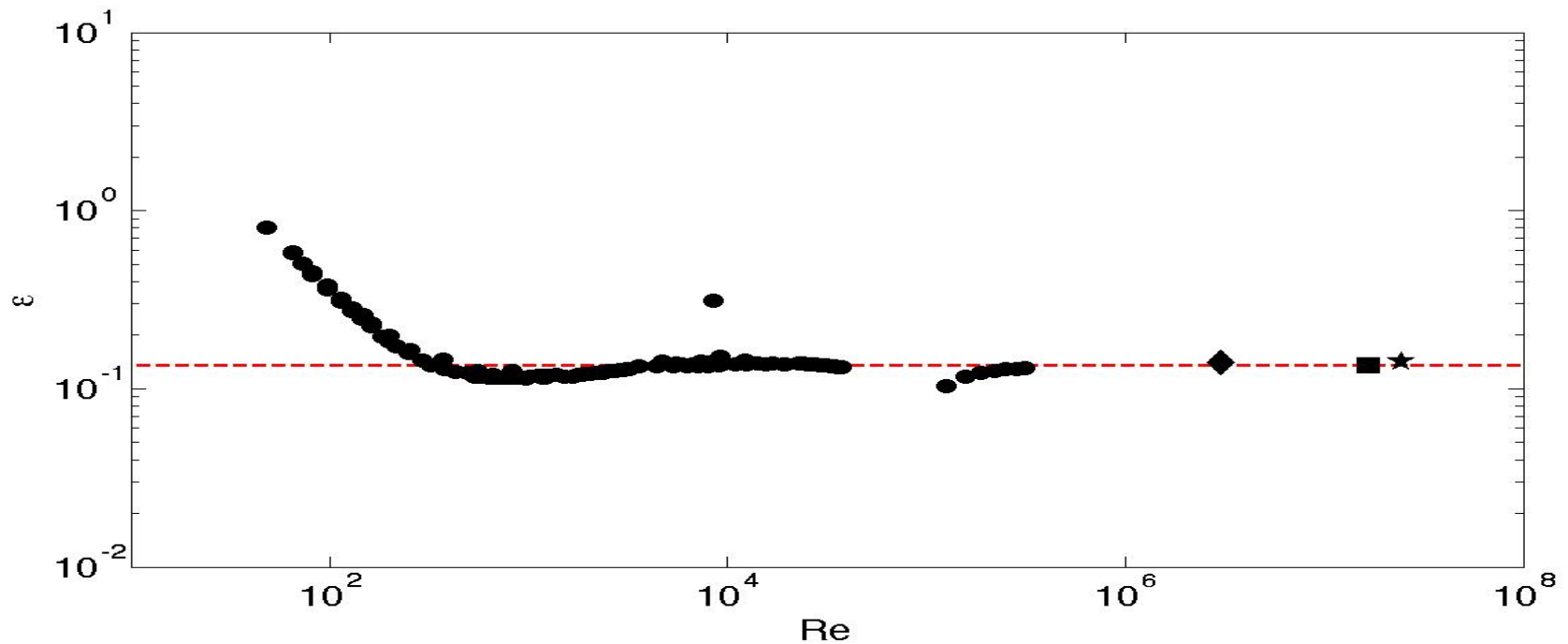
$$V_x(x, y, z), V_y(x, y, z), V_z(x, y, z)$$

Turbulent flow viewed in PTV



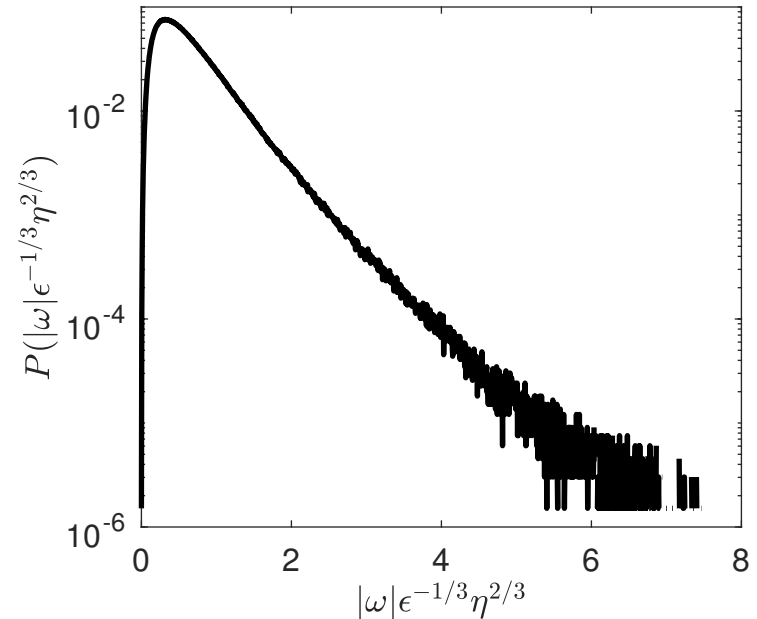
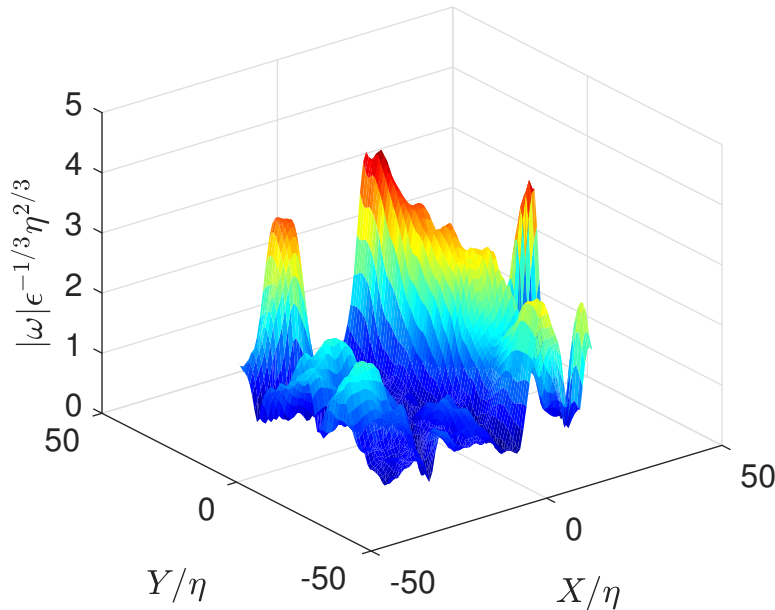
Time and space resolved Lagrangian trajectories

Energy dissipation in von Karman flow



Non-dimensional Energy dissipation per unit mass is constant at large Reynolds
Independent of viscosity?

Enstrophy blow-up



$$\epsilon = \nu \langle (\nabla u)^2 \rangle = \nu \langle \omega^2 \rangle \implies \langle \omega^2 \rangle \approx \frac{\epsilon}{\nu}$$

$$\lim_{\nu \rightarrow 0} \langle \omega^2 \rangle = \infty$$

***Building of very large gradients at small scale... Singularity?
How to measure them/quantify them in experiments?***

Singularities and Eulerian energy dissipation

$$-\mathcal{D} - \mathcal{D}_\nu$$

$$\frac{1}{2} \partial_t \mathbf{u}^2 + \operatorname{div} \left(\mathbf{u} \left(\frac{1}{2} \mathbf{u}^2 + p \right) - \nu \nabla \mathbf{u} \right) = D(u) - \nu \nabla \mathbf{u}^2$$

Inertial dissipation:

$$D(u) = \lim_{\ell \rightarrow 0} \frac{1}{4} \int_{r \leq \ell} d^3 r \nabla \phi_\ell(r) \cdot \delta u_r |\delta u_r|^2$$

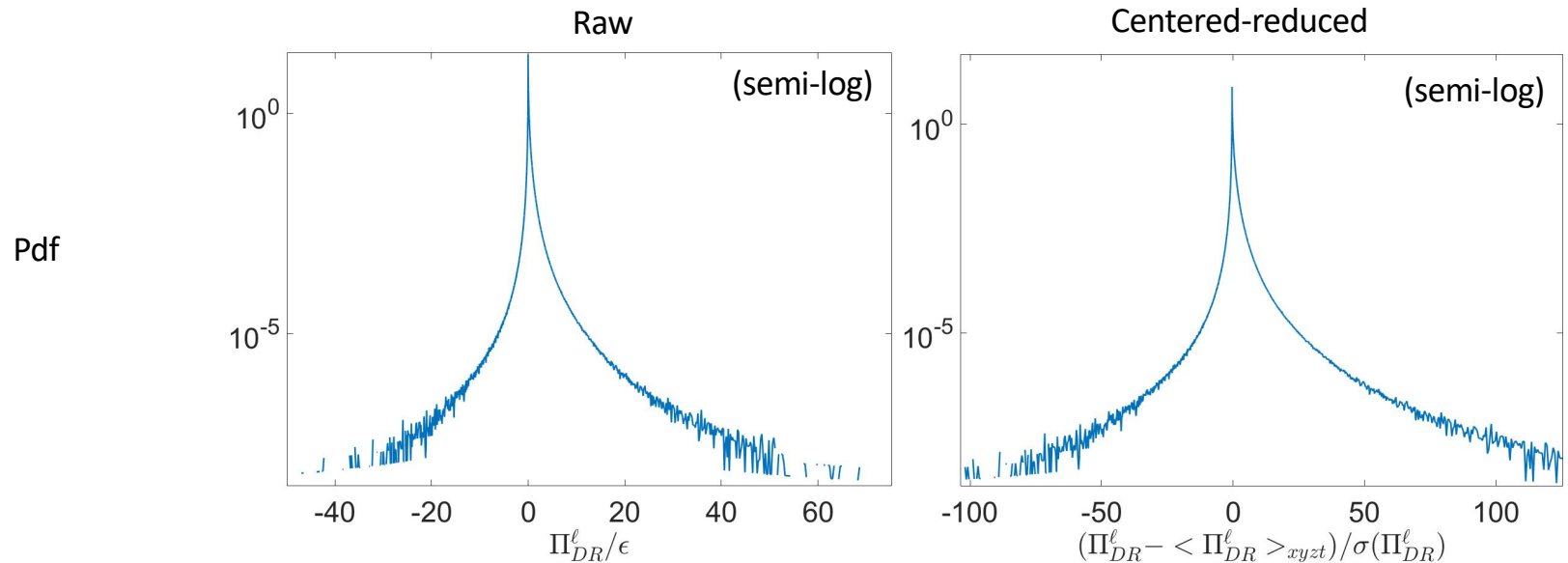
$$\delta u(\ell) \sim \ell^h \quad \text{In the limit of} \quad \ell \approx 0$$

$$D(u)[x] \propto \lim_{\ell \rightarrow 0} \ell^{3h-1}$$

If $h \leq 1/3 \rightarrow$ Dissipation through singularities

Duchon&Robert. Nonlinearity (2000)

Extreme events of the Duchon-Robert term

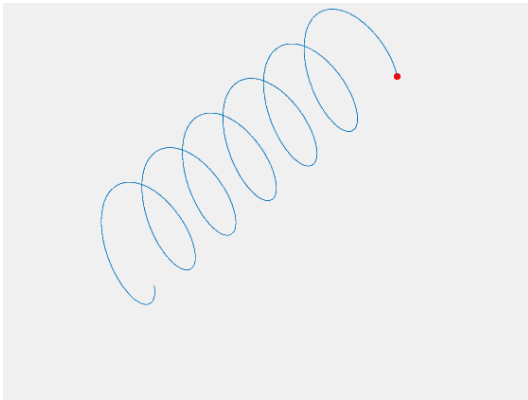


Largest $ \Pi_{DR}^\ell $	97ϵ	266σ
10^{th} largest $ \Pi_{DR}^\ell $	58ϵ	159σ
1000^{th} largest $ \Pi_{DR}^\ell $	14ϵ	38σ

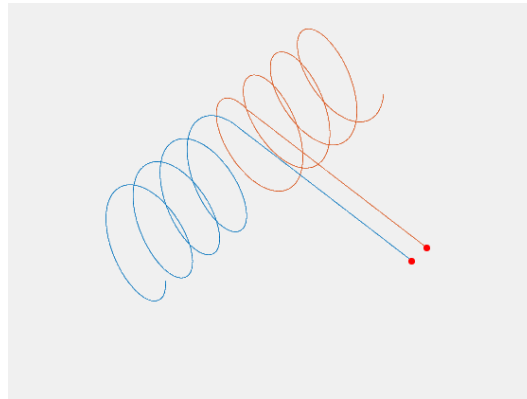
Direct observation of extreme events

- We observed three kinds of structures, based on the velocity streamlines :

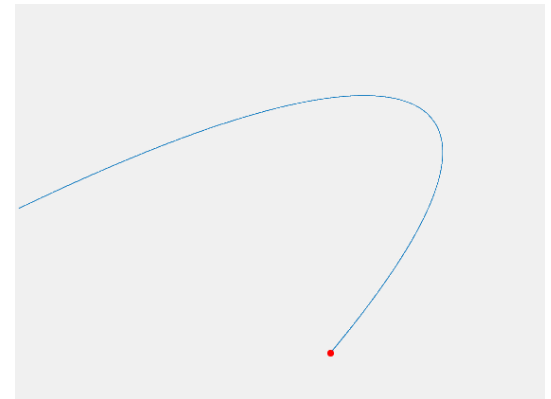
« Screw vortices »



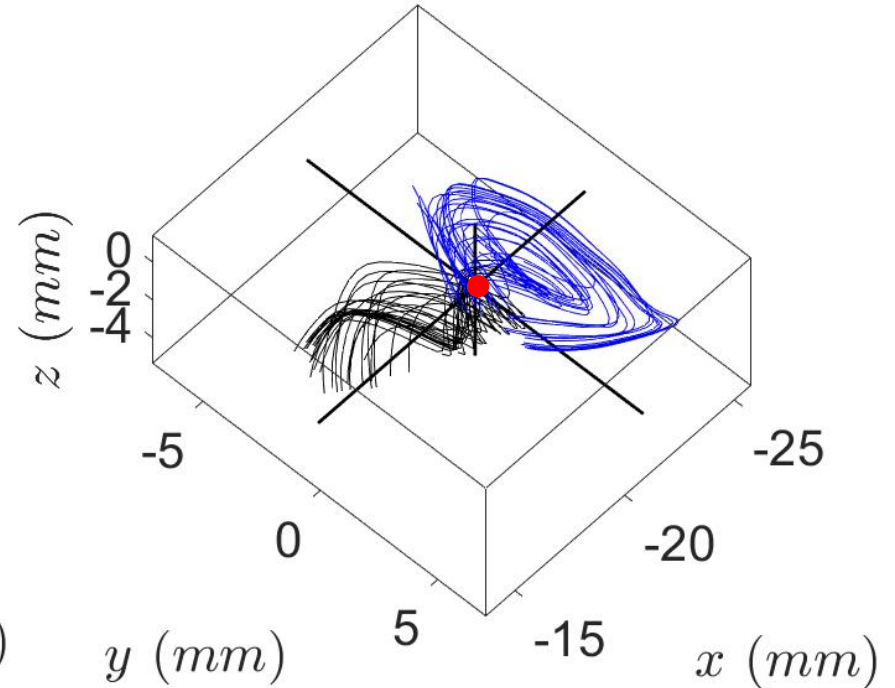
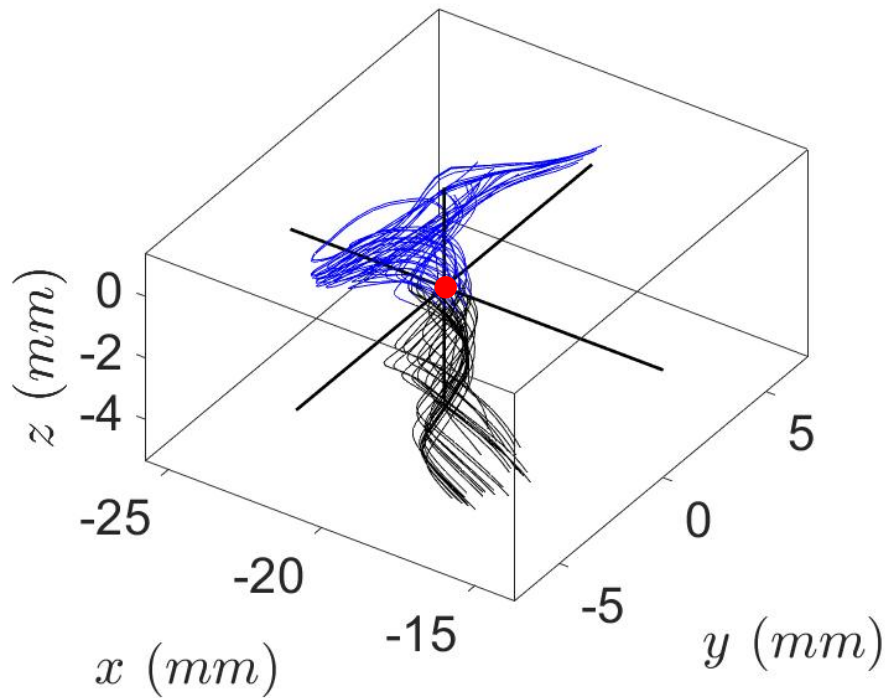
« Roll vortices »



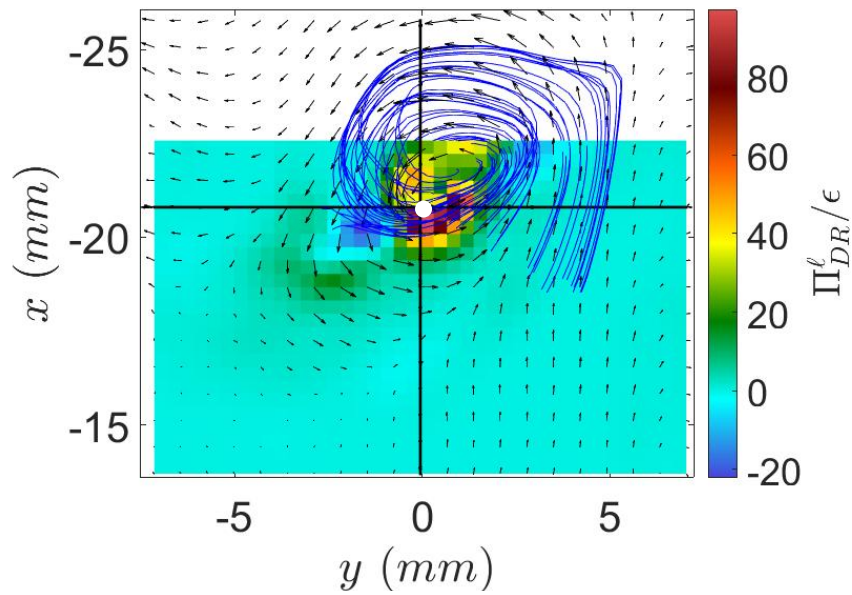
« U-turns »



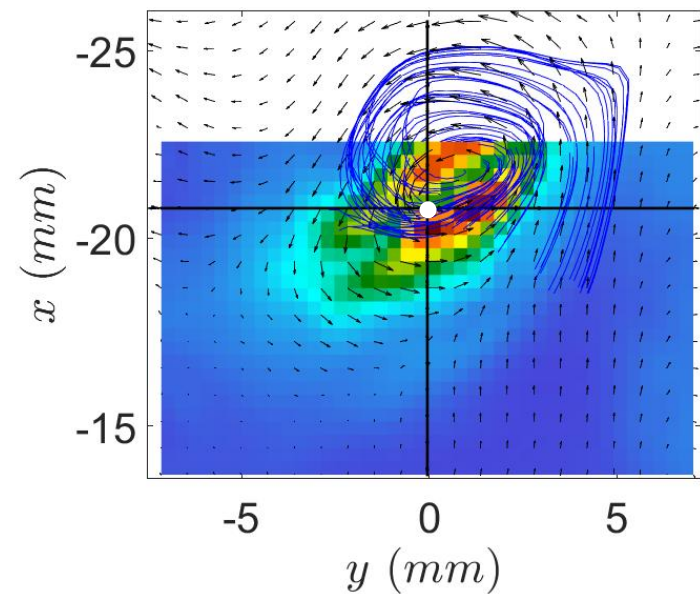
Screw vortices : velocity streamlines



Screw vortices : Π_{DR}^l and \mathcal{D}_l^v

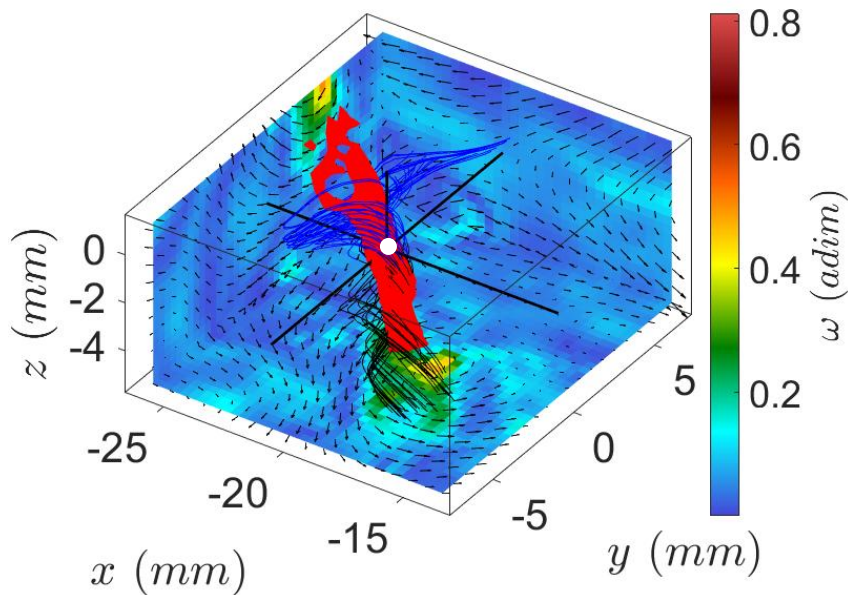


Π_{DR}^l

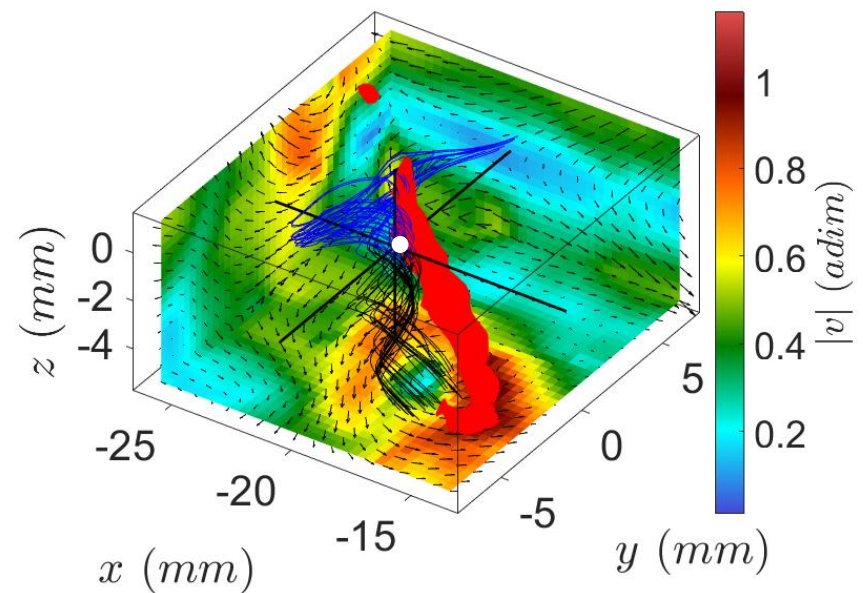


\mathcal{D}_l^v

Screw vortices : vorticity and velocity norms

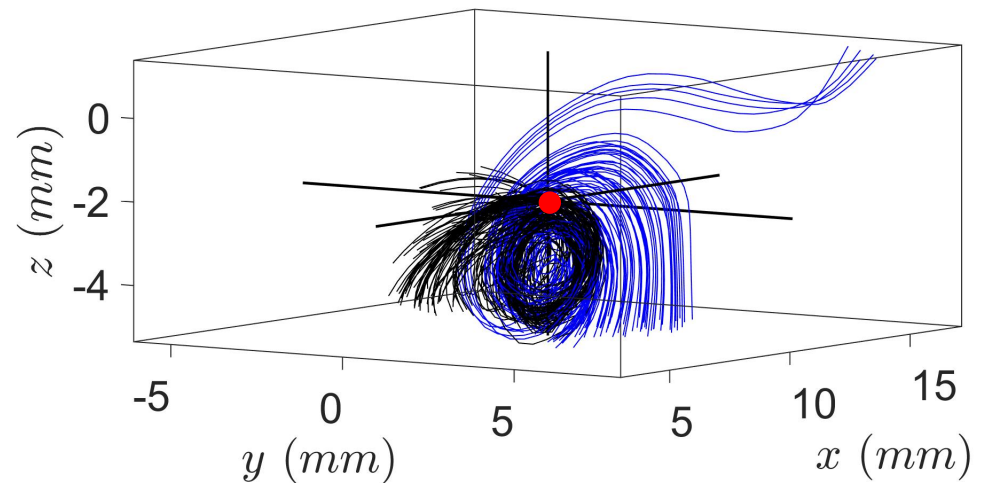
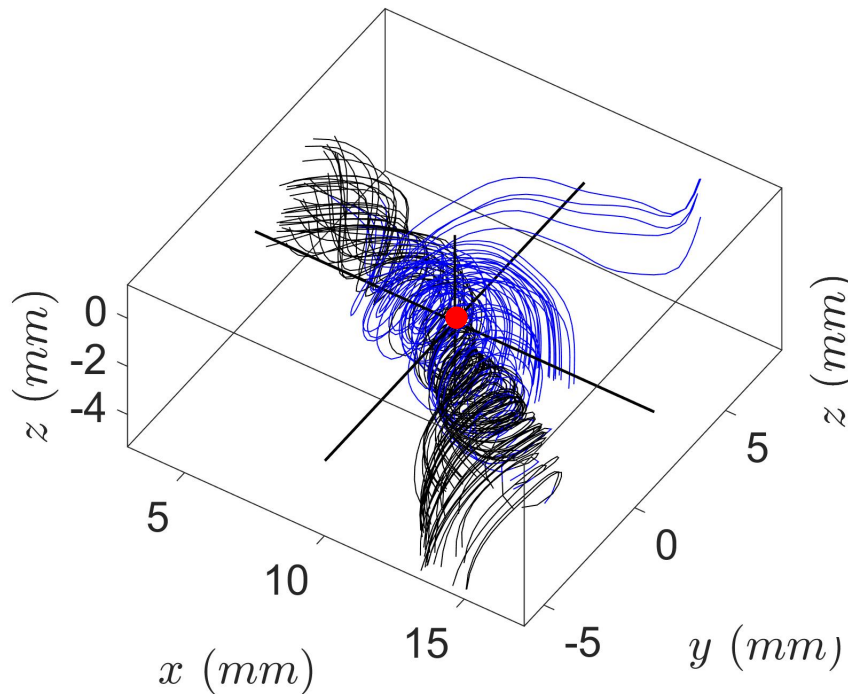


ω

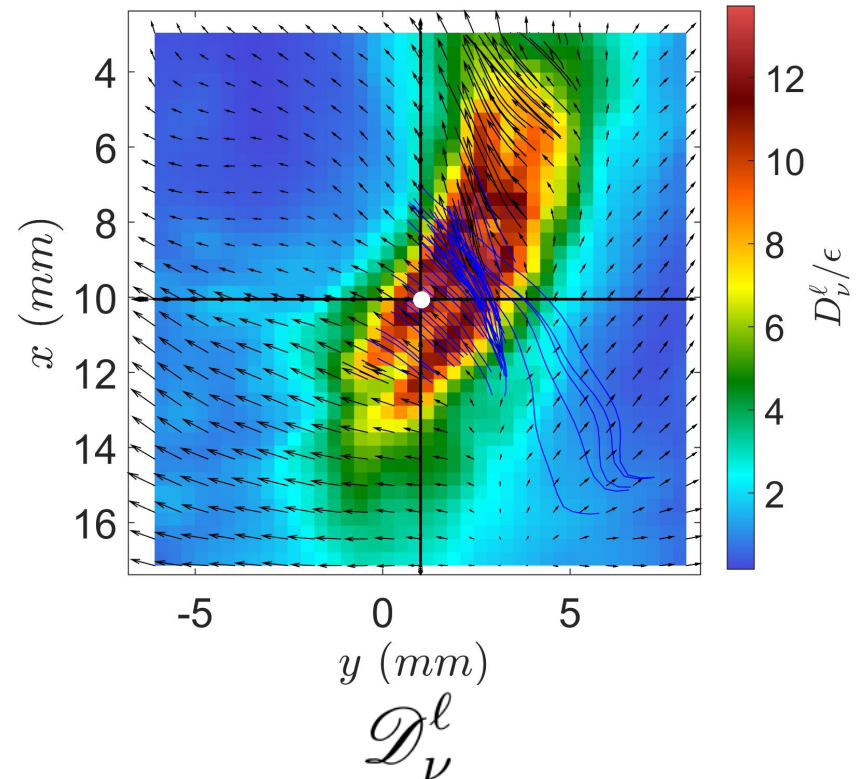
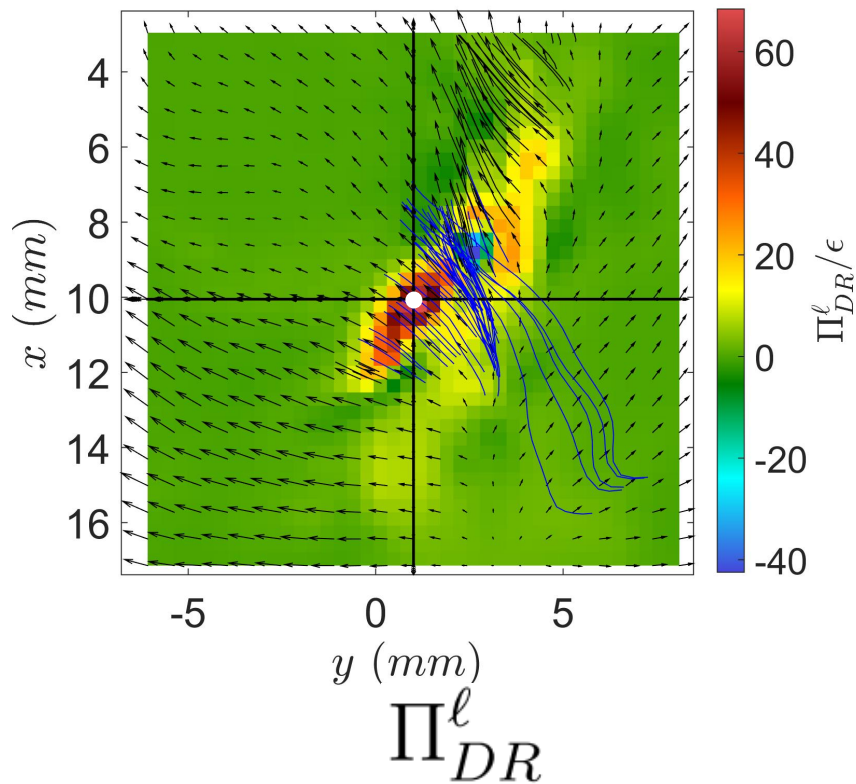


$|v|$

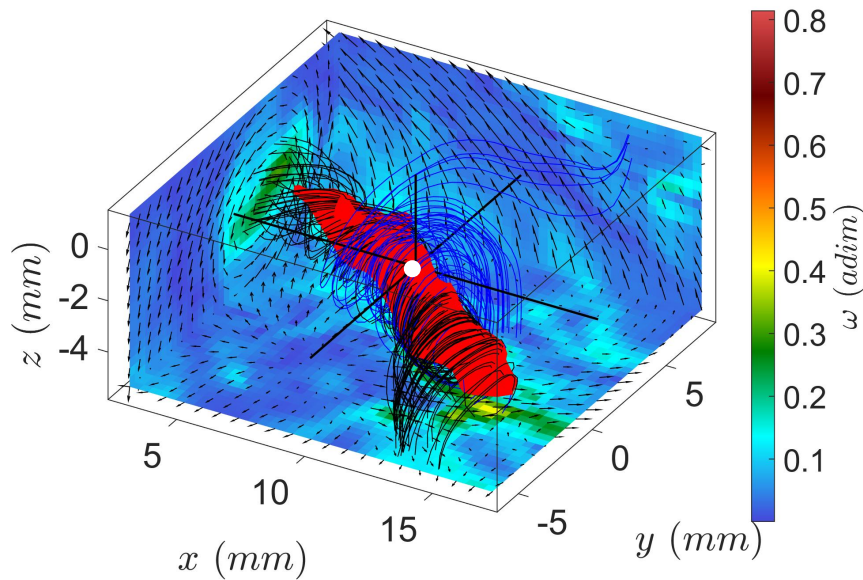
Roll vortices : velocity streamlines



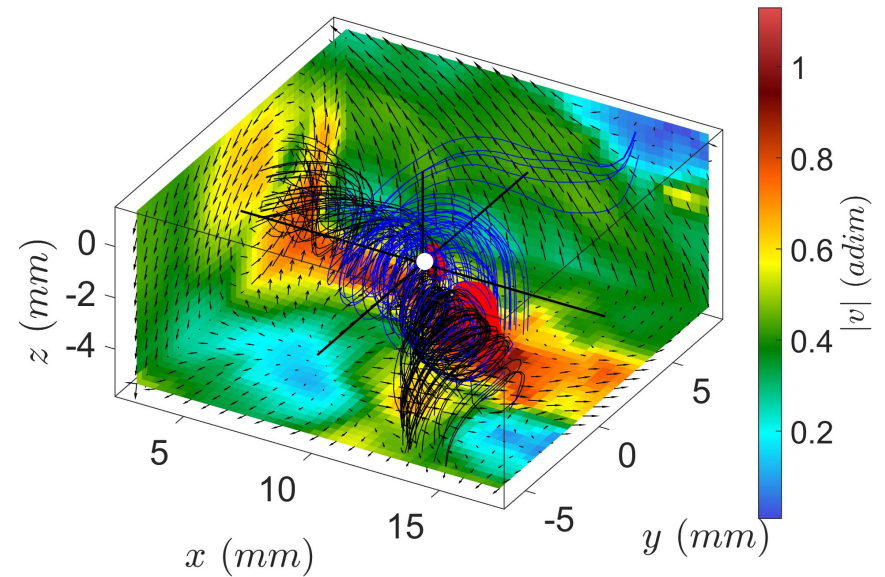
Roll vortices : Π_{DR}^l and \mathcal{D}_l^v



Roll vortices : vorticity and velocity norms

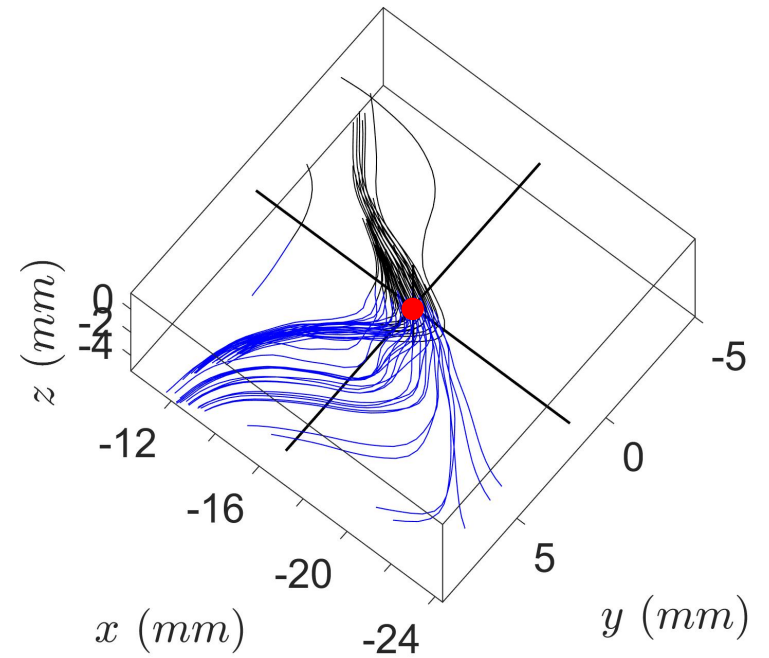
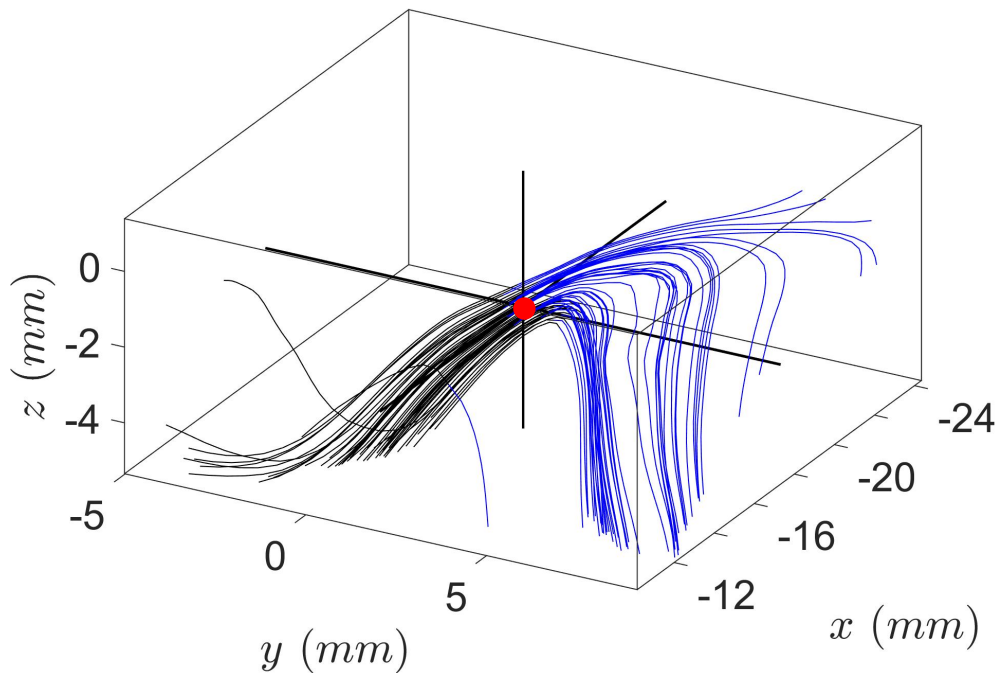


ω

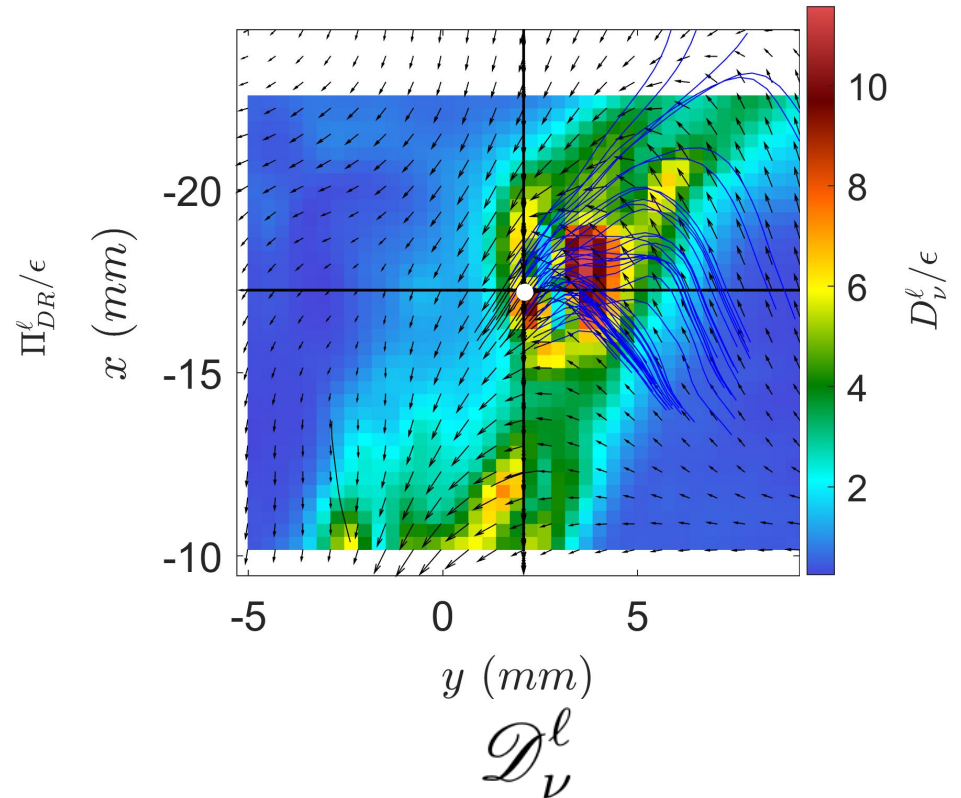
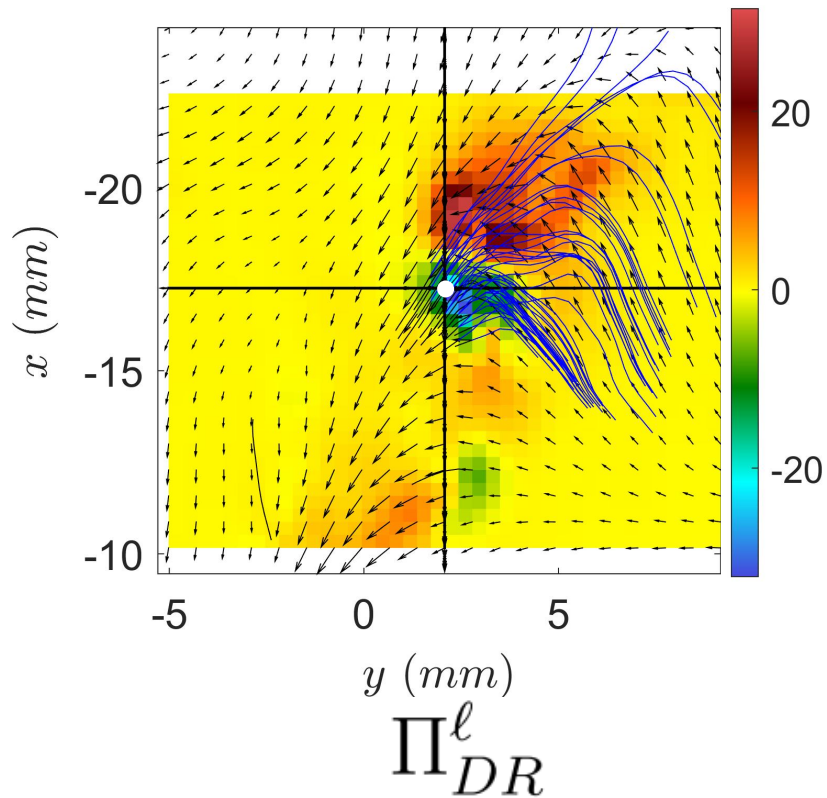


$|v|$

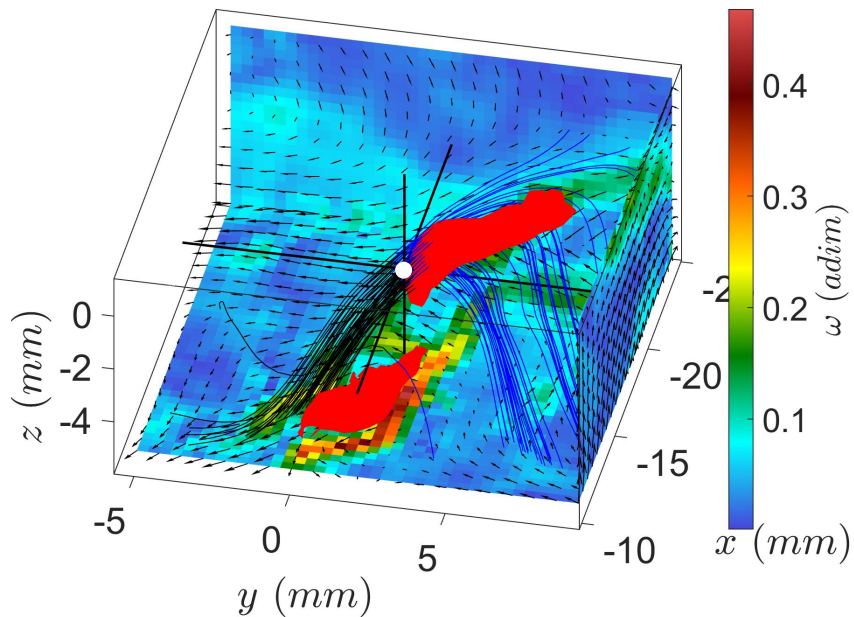
U-turns : velocity streamlines



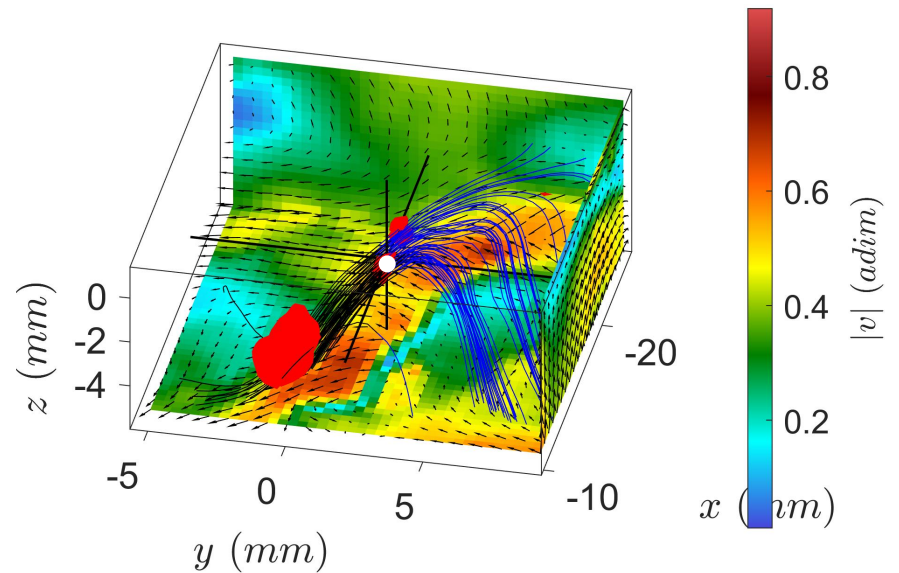
U-turns : Π_{DR}^l and \mathcal{D}_l^v



U-turns : vorticity and velocity norms



ω



$|v|$

Eulerian vs Lagrangian local energy dissipation

$$-\mathcal{D} - \mathcal{D}_\nu$$

$$\frac{1}{2} \partial_t \mathbf{u}^2 + \operatorname{div} \left(\mathbf{u} \left(\frac{1}{2} \mathbf{u}^2 + p \right) - \nu \nabla \mathbf{u} \right) = D(u) - \nu \nabla \mathbf{u}^2$$

Richardson law :

$$\left\langle (\delta X_r^{t \mp \tau})^2 \right\rangle \sim a_{\mp} \tau^3$$

Inertial dissipation:

$$D(u) = \lim_{\ell \rightarrow 0} \frac{1}{4} \int_{r \leq \ell} d^3 r \nabla \phi_\ell(r) \cdot \delta u_r |\delta u_r|^2$$

$$\delta u(\ell) \sim \ell^h \quad \text{In the limit of} \quad \ell \approx 0$$

$$D(u)[x] \propto \lim_{\ell \rightarrow 0} \ell^{3h-1}$$

If $h \leq 1/3 \rightarrow$ Dissipation through singularities

Irreversibility

$$a_+ \neq a_-$$

$$D_L = \lim_{r, \tau \rightarrow (0,0)} (\langle (\delta X_r^{t+\tau})^2 \rangle - \langle (\delta X_r^{t-\tau})^2 \rangle) / 6\tau^3$$

J. Jucha et al. PRL (2014)

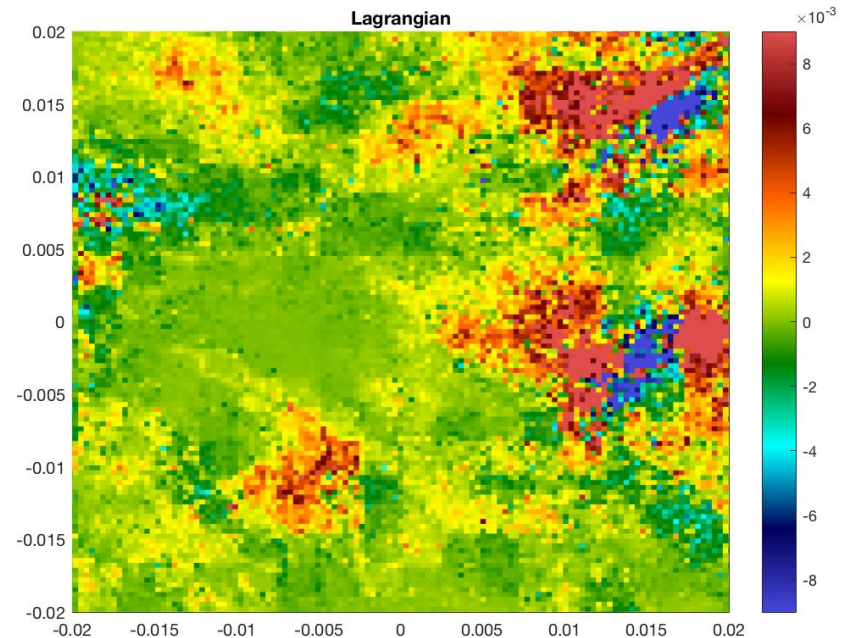
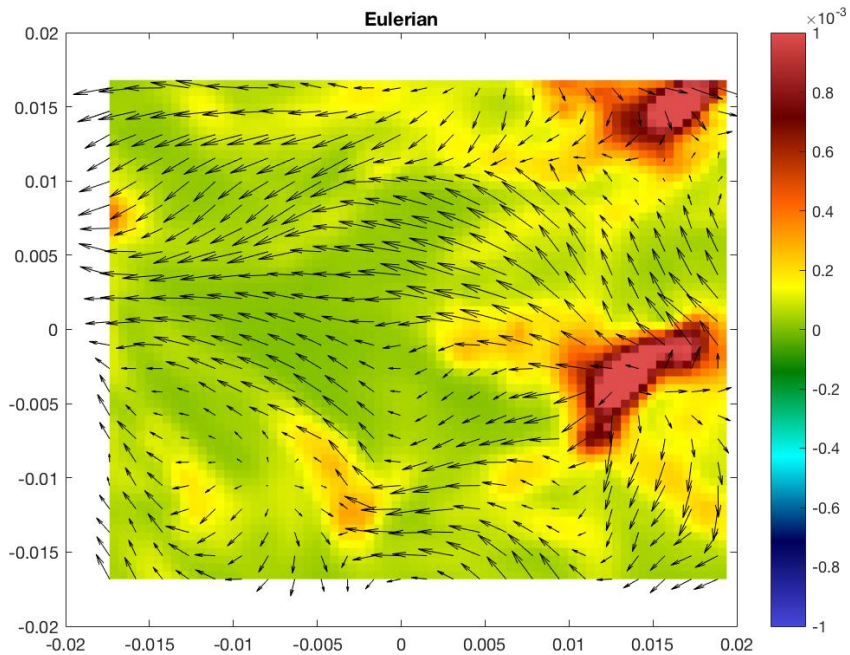
Duchon & Robert. Nonlinearity (2000)

Theory

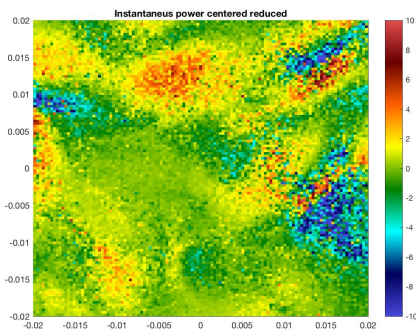
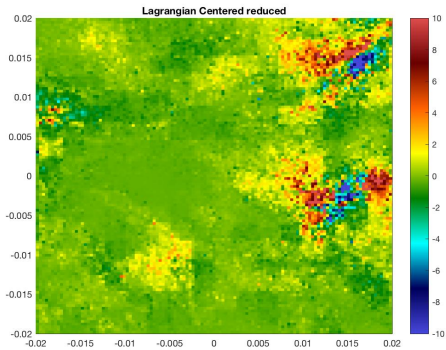
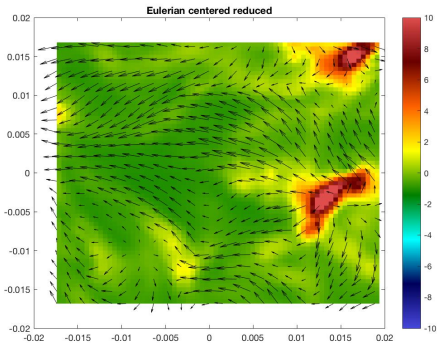
$$D_L = \varepsilon = \mathcal{D} + \mathcal{D}_\nu$$

Th. Drivas, 2019

Eulerian vs Lagrangian local dissipation



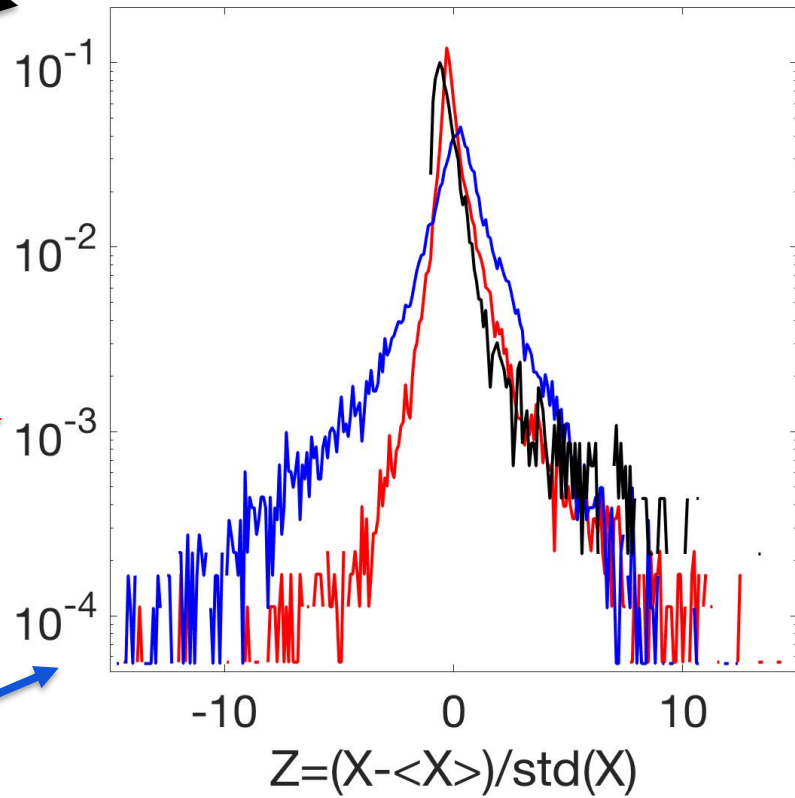
Eulerian vs Lagrangian local dissipation



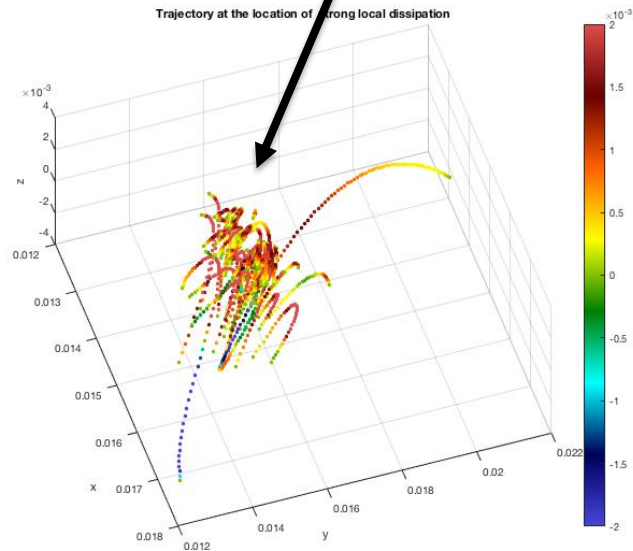
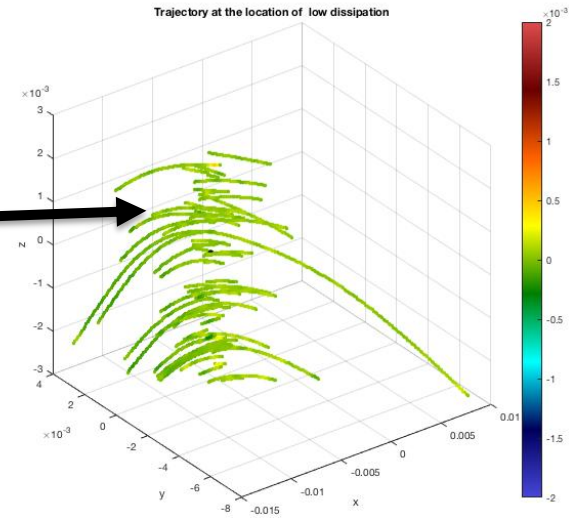
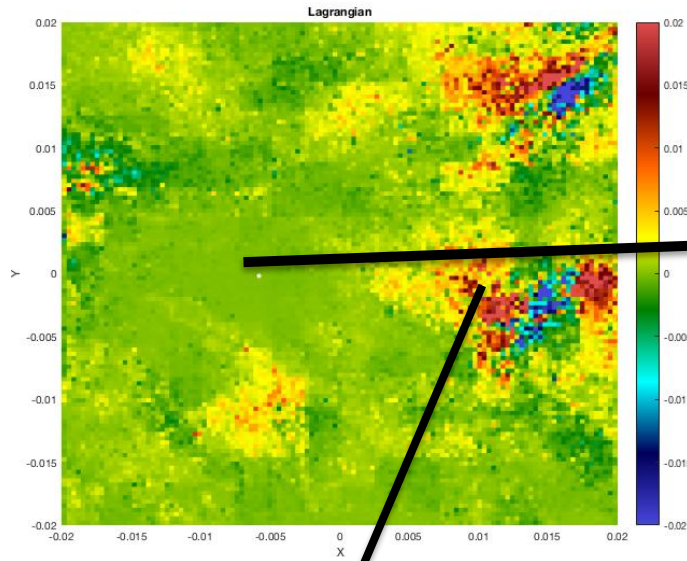
Instantaneous power different from local dissipation

Centered Reduced distributions

$P(Z)$

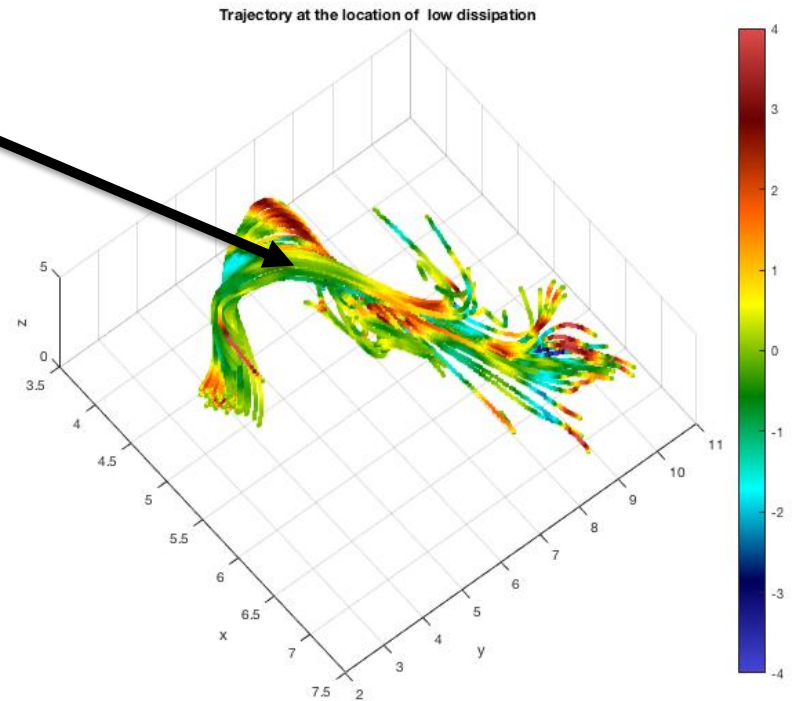
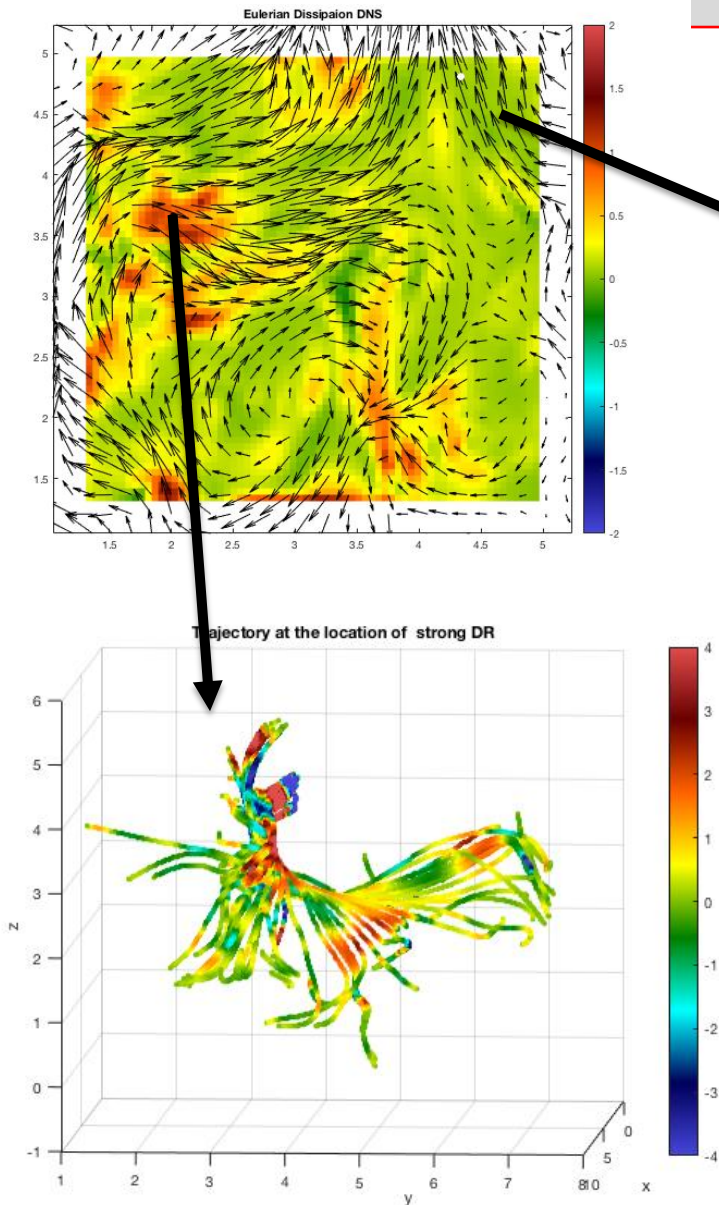


Geometry of trajectories at special points



Chaotic trajectories at high dissipation
Regular trajectories at low dissipation

Geometry of trajectories at special points (DNS)



Conclusion

Lagrangian and Eulerian irreversibility provides **complementary** view

Eulerian irreversibility points to **quasi singularities** of Eulerian Velocity Field

Lagrangian irreversibility points to **chaotic points** of Lagrangian trajectories

OPEN Questions: signatures of non-unicity?

Structure of quasi-singularities?

Dynamics of quasi-singularities?